

The Finite Element Method for the Analysis of Linear Systems



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Contents of Today's Lecture

• Formulation of structural elements

Beam and axis-symmetric shell elements

- Straight beam elements
- General curved beam elements
- Transition elements
- Axis-symmetric shell elements



Introduction:

Previously we considered beam elements in which we neglected the shear deformations (Bernoulli beams).

There are drawbacks of this approach:

- Its generalization to shells and plates is difficult, as it is difficult to satisfy continuity requirements to displacements and rotations (rotations are calculated from transverse displacements).

- Using flat elements to represent shells and plates necessitates a large number of elements to represent the geometry sufficiently accurately.





Introduction:

We will now consider beam, shell and plate elements where we take into account the shear deformations.

Furthermore we will apply iso-parametric formulations allowing for the accurate representation of general curved shell geometries.

Start is taken in considering the straight beam element – whereafter the concepts are applied for also shell and plate elements.



Straight beam elements: Neglecting shear effects





Straight beam elements: Including shear effects





Straight beam elements: Including shear effects









Two-dimensional beam elements:



For this beam (including shear strains) the principle of virtual work gives:

$$EI\int_{0}^{L} \left(\frac{d\beta}{dx}\right) \left(\frac{d\overline{\beta}}{dx}\right) dx + GAk\int_{0}^{L} \left(\frac{dw}{dx} - \beta\right) \left(\frac{d\overline{w}}{dx} - \overline{\beta}\right) dx = \int_{0}^{L} p\overline{w}dx + \int_{0}^{L} m\overline{\beta}dx$$









We can now assemble the element matrixes:

$$\mathbf{K} = EI \int_{-1}^{1} \mathbf{B}_{\beta}^{T} \mathbf{B}_{\beta}^{T} \det Jdr + GAk \int_{-1}^{1} (\mathbf{B}_{w}^{T} - \mathbf{H}_{\beta})^{T} (\mathbf{B}_{w}^{T} - \mathbf{H}_{\beta}) \det Jdr$$
$$\mathbf{R} = \int_{-1}^{1} \mathbf{H}_{W}^{T} p \det Jdr + \int_{-1}^{1} \mathbf{H}_{\beta}^{T} m \det Jdr$$



Two-dimensional beam elements:



The element we just derived is however only recommendable if we use 3 of 4 node elements and provided that the interior nodes are located at the midpoint or third points respectively – the reason for this is shear locking

To appreciate this phenomena we consider the total potential for a beam including the potential due to shear stresses



Two-dimensional beam elements:



The total potential can be written as:

$$\Pi = \frac{EI}{2} \int_{0}^{L} \left(\frac{d\beta}{dx}\right)^{2} dx + \frac{GAk}{2} \int_{0}^{L} \left(\frac{dw}{dx} - \beta\right)^{2} dx - \int_{0}^{L} pwdx + \int_{0}^{L} m\beta dx$$

The relative contribution from the bending and shear strain energies can be written as:

$$\tilde{\Pi} = \int_{0}^{L} \left(\frac{d\beta}{dx}\right)^{2} dx + \frac{GAk}{EI} \int_{0}^{L} \left(\frac{dw}{dx} - \beta\right)^{2} dx$$



Two-dimensional beam elements:



By study of this expression for different geometries it is evident that for small *h* the element must be able to represent zero shear strain conditions

$$\tilde{\Pi} = \int_{0}^{L} \left(\frac{d\beta}{dx}\right)^{2} dx + \frac{GAk}{EI} \int_{0}^{L} \left(\frac{dw}{dx} - \beta\right)^{2} dx$$
In such situations this term must be able to approach zero as well

For $h \rightarrow 0$ $I \rightarrow 0$ (but to the power of 3 - much faster!)

If elements are not able to represent zero shear strain conditions – we are dealing with shear locking – the element is too stiff!



Two-dimensional beam elements:

To study the phenomenon of shear locking lets consider a small example:







Two-dimensional beam elements:



However, we may impose the requirement that the shear strains are zero only at mid point:

$$\gamma = 0 = \frac{w_2}{L} - \frac{1}{2}\theta_2$$
 At midpoint !
$$w_2 = \frac{L}{2}\theta_2$$

This points to the idea that we could have assumed a constant shear strain equal to

 $\gamma = \frac{w_2}{L} - \frac{1}{2}\theta_2$

Mixed interpolation between displacements and rotations



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Two-dimensional beam elements:

Thus we get for 2, 3 and 4 node elements:





Two-dimensional beam elements:

Using mixed interpolation for the shear strains greatly improves the performance of the element:

In addition – the stiffness matrix can still be calculated exactly for what concerns the bending strain contributions (full integration) and then simply add the terms from the shear strains using the Gauss integration (mid point for a two node element)

This is sometimes referred to as reduced integration!







General curved beam elements:

The displacement components are found as:

$u(r,s,t) = \sum_{k=1}^{q} h_k^{\ell} u_k + \frac{t}{2} \sum_{k=1}^{q} a_k h_k V_{t,x}^k + \frac{s}{2} \sum_{k=1}^{q} b_k h_k V_{s,x}^k$ $u(r,s,t) = {}^{1}x - {}^{0}x$ $w(r, s, t) = {}^{1}z - {}^{0}z$ $w(r,s,t) = \sum_{k=1}^{q} h_k^{\ell} w_k + \frac{t}{2} \sum_{k=1}^{q} a_k h_k V_{t,z}^k + \frac{s}{2} \sum_{k=1}^{q} b_k h_k V_{s,z}^k$

Method of Finite Elements 1

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 ${}^{0}\mathbf{V}^{1}$

 $\begin{bmatrix} e & e \end{bmatrix}$



General curved beam elements:

The displacement interpolation matrix H is found by inserting:

$$\begin{vmatrix} \mathbf{V}_{t}^{k} = \mathbf{\theta}_{k} \times {}^{0} \mathbf{V}_{t}^{k} \\ \mathbf{V}_{s}^{k} = \mathbf{\theta}_{k} \times {}^{0} \mathbf{V}_{s}^{k} \end{vmatrix} \quad \mathbf{\theta}_{k} = \begin{bmatrix} \theta_{x}^{k} \\ \theta_{y}^{k} \\ \theta_{z}^{k} \end{bmatrix}$$

into:

$$u(r,s,t) = \sum_{k=1}^{q} h_{k}^{\ell} u_{k} + \frac{t}{2} \sum_{k=1}^{q} a_{k} h_{k} V_{t,x}^{k} + \frac{s}{2} \sum_{k=1}^{q} b_{k} h_{k} V_{s,x}^{k}$$
$$v(r,s,t) = \sum_{k=1}^{q} h_{k}^{\ell} v_{k} + \frac{t}{2} \sum_{k=1}^{q} a_{k} h_{k} V_{t,y}^{k} + \frac{s}{2} \sum_{k=1}^{q} b_{k} h_{k} V_{s,y}^{k}$$
$$w(r,s,t) = \sum_{k=1}^{q} h_{k}^{\ell} w_{k} + \frac{t}{2} \sum_{k=1}^{q} a_{k} h_{k} V_{t,z}^{k} + \frac{s}{2} \sum_{k=1}^{q} b_{k} h_{k} V_{s,z}^{k}$$





General curved beam elements:

For the beam only the longitudinal strain and the transversal shear strains are required:



and

$$\mathbf{B} = \begin{bmatrix} \mathbf{B}_1 & \cdots & \mathbf{B}_q \end{bmatrix}$$





General curved beam elements:

In order to evaluate the components of B we must take the derivatives of H following "V the usual procedure for iso-parametric elements:

$$\frac{\frac{\partial(u,v,w)}{\partial r}}{\frac{\partial(u,v,w)}{\partial s}} = \sum_{k=1}^{q} \begin{bmatrix} \frac{\partial h_{k}}{\partial r} \begin{bmatrix} 1 & (g)_{1i}^{k} & (g)_{2i}^{k} & (g)_{3i}^{k} \end{bmatrix}}{h_{k}} \begin{bmatrix} 0 & (\hat{g})_{1i}^{k} & (\hat{g})_{2i}^{k} & (\hat{g})_{3i}^{k} \end{bmatrix}} \begin{bmatrix} (u,v,w)_{k} \\ \theta_{x}^{k} \\ \theta_{y}^{k} \end{bmatrix}, \qquad i = 1 \text{ for } u$$

$$i = 2 \text{ for } v$$

$$i = 3 \text{ for } w$$

$$(\hat{\mathbf{g}})^{k} = \frac{b_{k}}{2} \begin{bmatrix} 0 & -{}^{0}V_{s,z}^{k} & {}^{0}V_{s,y}^{k} \\ -{}^{0}V_{s,z}^{k} & 0 & -{}^{0}V_{s,x}^{k} \\ -{}^{0}V_{s,y}^{k} & {}^{0}V_{s,x}^{k} & 0 \end{bmatrix}, \qquad (\overline{\mathbf{g}})^{k} = \frac{a_{k}}{2} \begin{bmatrix} 0 & -{}^{0}V_{t,z}^{k} & {}^{0}V_{t,y}^{k} \\ -{}^{0}V_{t,z}^{k} & 0 & -{}^{0}V_{t,x}^{k} \\ -{}^{0}V_{t,y}^{k} & {}^{0}V_{t,x}^{k} & 0 \end{bmatrix}$$

$$(g)_{ij}^{k} = s(\hat{g})_{ij}^{k} + t(\overline{g})_{ij}^{k}$$





General curved beam elements:

Now we may transform into the global coordinate system:

$$\begin{bmatrix} \frac{\partial u(u,v,w)}{\partial x} \\ \frac{\partial(u,v,w)}{\partial y} \\ \frac{\partial(u,v,w)}{\partial z} \end{bmatrix} = \sum_{k=1}^{q} \begin{bmatrix} J_{11}^{-1} \frac{\partial h_k}{\partial r} & (G1)_{i1}^k & (G2)_{i1}^k & (G3)_{i1}^k \\ J_{21}^{-1} \frac{\partial h_k}{\partial r} & (G1)_{i2}^k & (G2)_{i2}^k & (G3)_{i2}^k \\ J_{31}^{-1} \frac{\partial h_k}{\partial r} & (G1)_{i3}^k & (G2)_{i3}^k & (G3)_{i3}^k \end{bmatrix} \begin{bmatrix} u_k \\ \theta_x^k \\ \theta_y^k \\ \theta_z^k \end{bmatrix}, \qquad i = 1 \text{ for } u$$

Page 25 a_{γ} a_1 ${}^{0}\mathbf{V}_{t}^{1}$

v

where

$$(Gm)_{in}^{k} = \left[J_{n1}^{-1}(g)_{mi}^{k}\right] \frac{\partial h_{k}}{\partial r} + \left[J_{n2}^{-1}(\hat{g})_{mi}^{k} + J_{n3}^{-1}(\overline{g})_{mi}^{k}\right] h_{k}$$



General curved beam elements:

Finally all we need is the strain-stress matrix:



$$\begin{bmatrix} \tau_{\eta\eta} \\ \tau_{\eta\xi} \\ \tau_{\eta\zeta} \end{bmatrix} = \begin{bmatrix} E & 0 & 0 \\ 0 & Gk & 0 \\ 0 & 0 & Gk \end{bmatrix} \begin{bmatrix} \varepsilon_{\eta\eta} \\ \gamma_{\eta\xi} \\ \gamma_{\eta\zeta} \end{bmatrix}$$

where *k* is the shear correction factor







General curved beam elements: Example We consider the following simple beam:



Page 27 Page 27 r, η r, η r,

 a_1





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Formulation of Structural Elements General curved beam elements: Example We consider the following simple beam: $V_{t}^{k} = \theta_{k} \times {}^{0}V_{t}^{k} = det \begin{bmatrix} e_{x} & e_{y} & e_{z} \\ \theta_{x}^{k} & \theta_{y}^{k} & \theta_{z}^{k} \end{bmatrix} = \theta_{y}^{k}e_{x} - \theta_{x}^{k}e_{y}$

$$\mathbf{V}_{s}^{k} = \mathbf{\theta}_{k} \times {}^{0}\mathbf{V}_{s}^{k} = \det \begin{bmatrix} e_{x} & e_{y} & e_{z} \\ \theta_{x}^{k} & \theta_{y}^{k} & \theta_{z}^{k} \\ 0 & 1 & 0 \end{bmatrix} = \theta_{z}^{k}e_{x} - \theta_{x}^{k}e_{z}$$

as we consider a 2-dimensional beam there is:

$$\theta_x^k = \theta_y^k = 0 \Longrightarrow \mathbf{V}_t^k = 0, \quad \mathbf{V}_s^k = \theta_z^k e_x$$



General curved beam elements: Example

Furthermore we have that only displacements in the y-direction are possible



$$u(r,s,t) = \sum_{k=1}^{q} h_{k}^{\ell} u_{k} + \frac{t}{2} \sum_{k=1}^{q} a_{k} h_{k} V_{t,x}^{k} + \frac{s}{2} \sum_{k=1}^{q} b_{k} h_{k} V_{s,x}^{k}$$

$$v(r,s,t) = \sum_{k=1}^{q} h_{k}^{\ell} v_{k} + \frac{t}{2} \sum_{k=1}^{q} a_{k} h_{k} V_{t,y}^{k} + \frac{s}{2} \sum_{k=1}^{q} b_{k} h_{k} V_{s,y}^{k}$$

$$w(r,s,t) = \sum_{k=1}^{q} h_{k}^{\ell} w_{k} + \frac{t}{2} \sum_{k=1}^{q} a_{k} h_{k} V_{t,z}^{k} + \frac{s}{2} \sum_{k=1}^{q} b_{k} h_{k} V_{s,z}^{k}$$

$$v(r) = \sum_{k=1}^{q} h_{k} v_{k}$$



Formulation of Structural Elements 2L3 3 **General curved beam elements: Example** Now we proceed to take the derivatives to develop the B matrix $\begin{vmatrix} \frac{\partial u}{\partial r} \\ \frac{\partial u}{\partial z} \end{vmatrix} = \sum_{k=1}^{3} \begin{vmatrix} \frac{\partial h_{k}}{\partial r} & \left[1 & (g)_{31}^{k} \right] \\ h_{k} & \left[0 & (\overline{g})_{31}^{k} \right] \end{vmatrix} \begin{vmatrix} u_{k} \\ \theta_{z}^{k} \end{vmatrix}, \quad u_{k} = 0$ $(\overline{\mathbf{g}})^{k} = \frac{a_{k}}{2} \begin{vmatrix} 0 & -{}^{0}V_{t,z}^{k} & {}^{0}V_{t,y}^{k} \\ -{}^{0}V_{t,z}^{k} & 0 & -{}^{0}V_{t,x}^{k} \\ -{}^{0}V^{k} & {}^{0}V^{k} & 0 \end{vmatrix}$][$\begin{vmatrix} \frac{\partial u}{\partial r} \\ \frac{\partial u}{\partial z} \end{vmatrix} = \sum_{k=1}^{3} \begin{bmatrix} \frac{\partial h_k}{\partial r} (g)_{31}^k \\ h_k (\overline{g})_{31}^k \end{bmatrix} \theta_z^k$ $(g)_{ii}^k = s(\overline{g})_{ij}^k$ $(\overline{g})_{31}^{k} = -\frac{h}{2} {}^{0}V_{t,y}^{k} = -\frac{h}{2}$ $(g)_{31}^{k} = s(\hat{g})_{31}^{k} = -\frac{h}{2}s^{0}V_{t,y}^{k} = -\frac{h}{2}s$ Method of Finite Elements 1



General curved beam elements: Example ² Now we proceed to take the derivatives to develop the B matrix





General curved beam elements: Example ² The next step is to transform into global coordinates



$$\frac{\partial}{\partial \mathbf{x}} = \mathbf{J}^{-1} \frac{\partial}{\partial \mathbf{r}}$$

$$\mathbf{J} = \begin{bmatrix} \frac{\partial x}{\partial r} & \frac{\partial y}{\partial r} \\ \frac{\partial x}{\partial s} & \frac{\partial y}{\partial s} \end{bmatrix}$$

$$\mathbf{J} = \begin{bmatrix} \frac{L}{2} & 0 \\ 0 & \frac{h}{2} \end{bmatrix} \longrightarrow \mathbf{J}^{-1} = 2 \begin{bmatrix} \frac{1}{L} & 0 \\ 0 & \frac{1}{h} \end{bmatrix}$$

$$\mathbf{x} = \frac{L}{2}(1+r)$$

$$\mathbf{y} = \frac{h}{2}(1+s)$$



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Formulation of Structural Elements

General curved beam elements: Example The next step is to transform into global coordinates





Now we can form the element matrixes as usual

 $\frac{2L}{3}$



General curved beam elements: Example The stiffness matrix is found from:

$$\mathbf{K} = \int_{-1}^{1} \int_{-1}^{1} \left[\varepsilon_{\eta\eta} \quad \gamma_{\eta\zeta} \right] \begin{bmatrix} E & 0 \\ 0 & Gk \end{bmatrix} \begin{bmatrix} \varepsilon_{\eta\eta} \\ \gamma_{\eta\zeta} \end{bmatrix} \det \mathbf{J} \, ds dr$$





Transition element

It is interesting to notice that the structural elements can be derived from the continuum elements by imposition of the deformation assumptions (degeneration)

This result facilitates the development of transition element between continuum elements and structural elements



x, u

Formulation of Structural Elements

Axi-symmetric shell element

This element can be constructed from the iso-parametric beam element (the 2-dimensional case; movements in the x,y-plane) $v_{i}v_{j}$

The stress-strain matrix must then be ammended by a row corresponding to the hoop strain – and the stress-strain law must include a term coupling between the hoop and the *r*-direction and ensuring that the stress in the *s*-direction is equal to zero

K.J. Bathe (Example 5.9 and Exercise 5.41)