

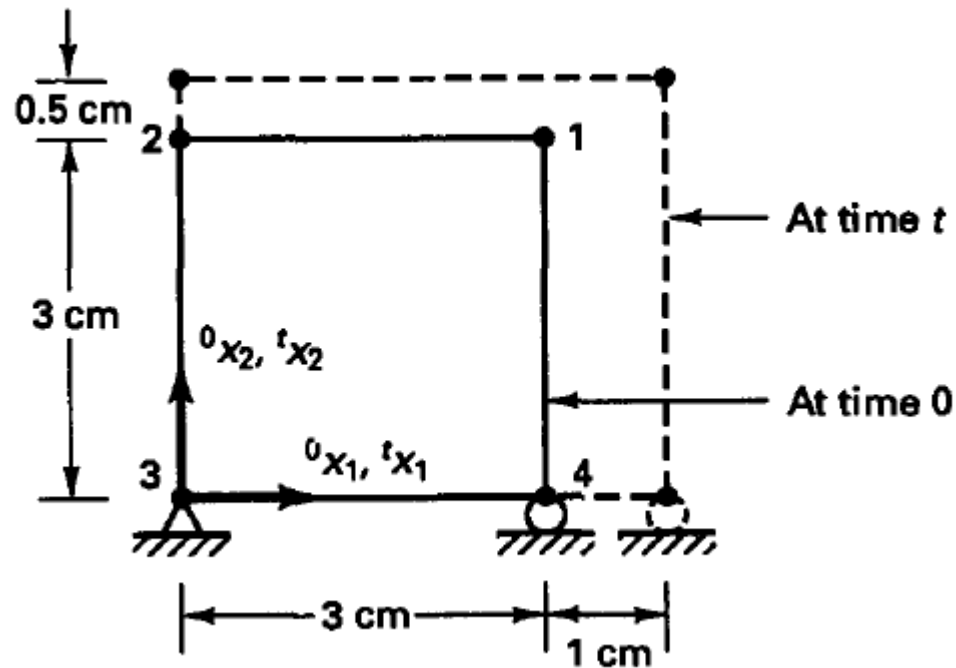
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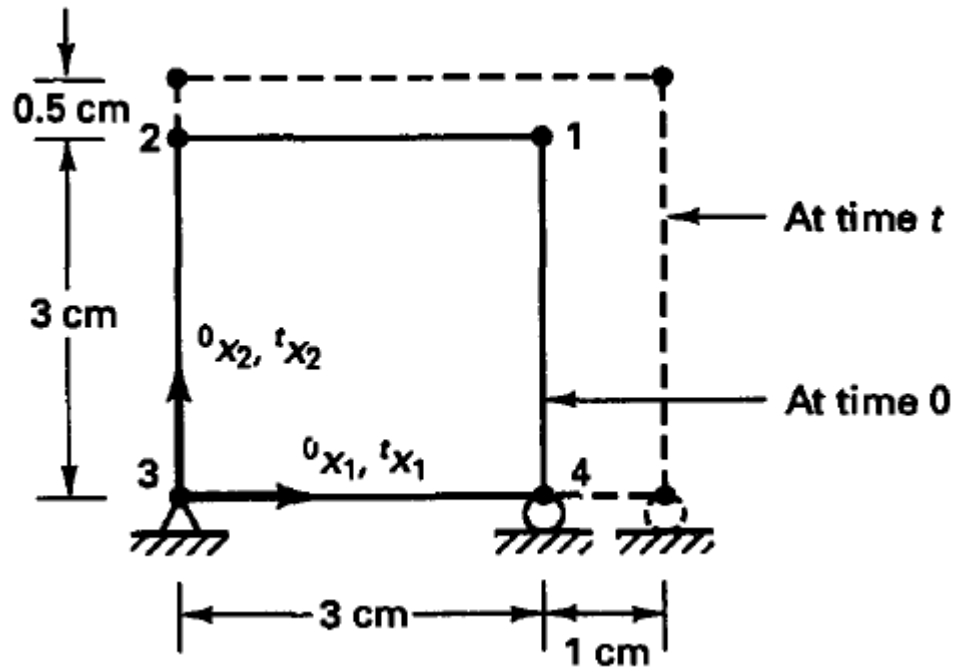
Finities Element Procedure, Klaus-Jürgen Bathe

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Example 6.18

Four-node plane strain element in large displacement/large strain conditions





Quest: Establish the matrices ${}^0\mathbf{B}_{L0}$, ${}^0\mathbf{B}_{L1}$, ${}^0\mathbf{B}_{NL}$
 Corresponding to TL using Table 6.5

Reminder

- TL – Total Lagrangian formulation

$$({}^t_0\mathbf{K}_L + {}^t_0\mathbf{K}_{NL})\mathbf{U} = {}^{t+\Delta t}\mathbf{R} - {}^t_0\mathbf{F}$$

$${}^t_0\mathbf{K}_L \hat{\mathbf{u}} = \left(\int_{{}_0V} {}^t_0\mathbf{B}_L^T {}_0\mathbf{C} {}^t_0\mathbf{B}_L d^0V \right) \hat{\mathbf{u}}$$

$${}^t_0\mathbf{K}_{NL} \hat{\mathbf{u}} = \left(\int_{{}_0V} {}^t_0\mathbf{B}_{NL}^T {}^t_0\mathbf{S} {}^t_0\mathbf{B}_{NL} d^0V \right) \hat{\mathbf{u}}$$

$${}^t_0\mathbf{F} = \int_{{}_0V} {}^t_0\mathbf{B}_L^T {}^t_0\hat{\mathbf{S}} d^0V$$

To get: \mathbf{U} = vector of increments in the nodal point displacements

We need to know: ${}_0\mathbf{C}$ = incremental stress-strain material property matrices

${}^t_0\mathbf{S}$, ${}^t_0\hat{\mathbf{S}}$ = matrix and vector of second Piola-Kirchhoff stresses

${}^t_0\mathbf{B}_L$ = linear strain-displacement transformation matrices

${}^t_0\mathbf{B}_{NL}$ = nonlinear strain-displacement transformation matrices

Table 6.5

$${}^0\mathbf{B}_L = {}^0\mathbf{B}_{L0} + {}^0\mathbf{B}_{L1}$$

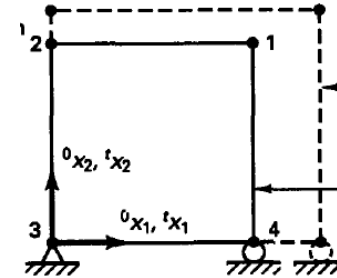
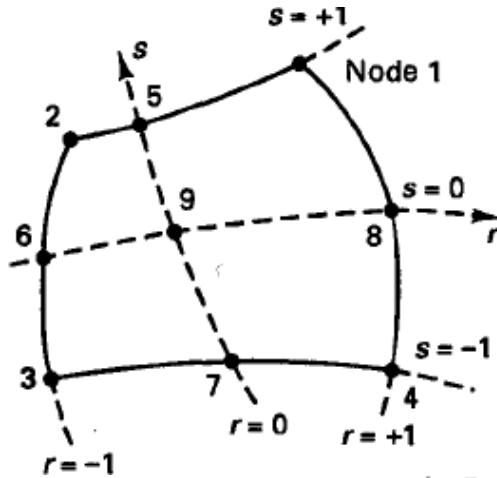
= linear strain-displacement transformation matrices

$${}^0\mathbf{B}_{L0} = \begin{bmatrix} {}^0h_{1,1} & 0 & {}^0h_{2,1} & 0 & {}^0h_{3,1} & 0 & \cdots & {}^0h_{N,1} & 0 \\ 0 & {}^0h_{1,2} & 0 & {}^0h_{2,2} & 0 & {}^0h_{3,2} & \cdots & 0 & {}^0h_{N,2} \\ {}^0h_{1,2} & {}^0h_{1,1} & {}^0h_{2,2} & {}^0h_{2,1} & {}^0h_{3,2} & {}^0h_{3,1} & \cdots & {}^0h_{N,2} & {}^0h_{N,1} \\ \frac{h_1}{{}^0\bar{x}_1} & 0 & \frac{h_2}{{}^0\bar{x}_1} & 0 & \frac{h_3}{{}^0\bar{x}_1} & 0 & \cdots & \frac{h_N}{{}^0\bar{x}_1} & 0 \end{bmatrix}$$

Only for 3D axisymmetric elements

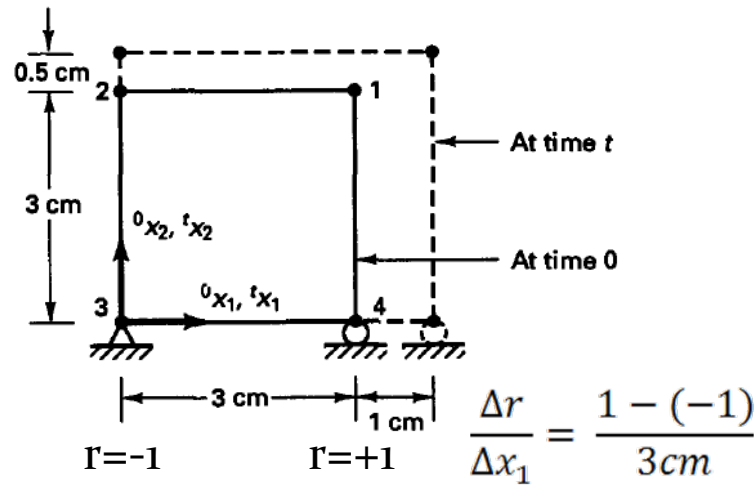
$$\text{where } {}^0h_{k,j} = \frac{\partial h_k}{\partial {}^0x_j};$$

Quadrilateral Elements



Include only if node i is defined

	$i = 5$	$i = 6$	$i = 7$	$i = 8$	$i = 9$
$h_1 = \frac{1}{4}(1+r)(1+s)$	$-\frac{1}{2}h_5$			$-\frac{1}{2}h_8$	$-\frac{1}{4}h_9$
$h_2 = \frac{1}{4}(1-r)(1+s)$	$-\frac{1}{2}h_5$	$-\frac{1}{2}h_6$			$-\frac{1}{4}h_9$
$h_3 = \frac{1}{4}(1-r)(1-s)$		$-\frac{1}{2}h_6$	$-\frac{1}{2}h_7$		$-\frac{1}{4}h_9$
$h_4 = \frac{1}{4}(1+r)(1-s)$			$-\frac{1}{2}h_7$	$-\frac{1}{2}h_8$	$-\frac{1}{4}h_9$
$h_5 = \frac{1}{2}(1-r^2)(1+s)$					$-\frac{1}{2}h_9$
$h_6 = \frac{1}{2}(1-s^2)(1-r)$					$-\frac{1}{2}h_9$
$h_7 = \frac{1}{2}(1-r^2)(1-s)$					$-\frac{1}{2}h_9$
$h_8 = \frac{1}{2}(1-s^2)(1+r)$					$-\frac{1}{2}h_9$
$h_9 = (1-r^2)(1-s^2)$					



$$h_1 = \frac{1}{4}(1+r)(1+s) \rightarrow {}_0h_{1,1} = \frac{dh_1}{dx_1} = \frac{dh_1}{dr} * \frac{dr}{dx_1} = \frac{1}{4}(1+s) * \frac{2}{3} = \frac{1}{6}(1+s)$$

$${}^t\mathbf{B}_{LO} = \frac{1}{6} \begin{bmatrix} (1+s) & 0 & -(1+s) & 0 & -(1-s) & 0 & (1-s) & 0 \\ 0 & (1+r) & 0 & (1-r) & 0 & -(1-r) & 0 & -(1+r) \\ (1+r) & (1+s) & (1-r) & -(1+s) & -(1-r) & -(1-s) & -(1+r) & (1-s) \end{bmatrix}$$

Table 6.5

$${}^t\mathbf{B}_{L1} = \begin{bmatrix} l_{11} {}^0h_{1,1} & l_{21} {}^0h_{1,1} & l_{11} {}^0h_{2,1} & l_{21} {}^0h_{2,1} & \dots & l_{11} {}^0h_{N,1} & l_{21} {}^0h_{N,1} \\ l_{12} {}^0h_{1,2} & l_{22} {}^0h_{1,2} & l_{12} {}^0h_{2,2} & l_{22} {}^0h_{2,2} & \dots & l_{12} {}^0h_{N,2} & l_{22} {}^0h_{N,2} \\ (l_{11} {}^0h_{1,2} + l_{12} {}^0h_{1,1}) & (l_{21} {}^0h_{1,2} + l_{22} {}^0h_{1,1}) & (l_{11} {}^0h_{2,2} + l_{12} {}^0h_{2,1}) & (l_{21} {}^0h_{2,2} + l_{22} {}^0h_{2,1}) & \dots & (l_{11} {}^0h_{N,2} + l_{12} {}^0h_{N,1}) & (l_{21} {}^0h_{N,2} + l_{22} {}^0h_{N,1}) \\ l_{33} \frac{h_1}{{}^0x_1} & 0 & l_{33} \frac{h_2}{{}^0x_1} & 0 & \dots & l_{33} \frac{h_N}{{}^0x_1} & 0 \end{bmatrix}$$

where $l_{11} = \sum_{k=1}^N {}^0h_{k,1} \text{ } {}^t u_1^k$;

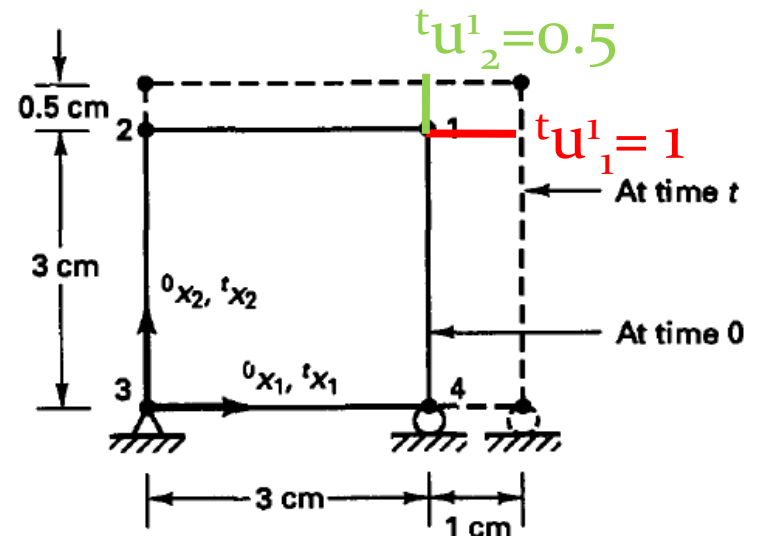
$$u_j^k = {}^{t+\Delta t} u_j^k - {}^t u_j^k;$$

$${}^t u_1^1 = 1; \quad {}^t u_2^1 = 0.5$$

$${}^t u_1^2 = 0; \quad {}^t u_2^2 = 0.5$$

$${}^t u_1^3 = 0; \quad {}^t u_2^3 = 0$$

$${}^t u_1^4 = 1; \quad {}^t u_2^4 = 0$$



$$l_{11} = \sum_{k=1}^4 {}_0h_{k,1} {}^t u_1^k = ({}_0h_{1,1} * {}^t u_1^1 + {}_0h_{4,1} * {}^t u_1^4) = \frac{1}{3}$$

$$h_1 = \frac{1}{4}(1+r)(1+s) \rightarrow {}_0h_{1,1} = \frac{1}{6}(1+s)$$

$$h_4 = \frac{1}{4}(1+r)(1-s) \rightarrow {}_0h_{4,1} = \frac{1}{6}(1-s)$$

$${}^t u_1^1 = 1;$$

$${}^t u_1^2 = 0;$$

$${}^t u_1^3 = 0;$$

$${}^t u_1^4 = 1;$$

$${}^t \mathbf{B}_{L1} = \begin{bmatrix} l_{11} {}_0h_{1,1} & \dots \\ l_{12} {}_0h_{1,2} & \dots \\ (l_{11} {}_0h_{1,2} + l_{12} {}_0h_{1,1}) & \dots \end{bmatrix}$$

$${}^t \mathbf{B}_{L1} =$$

$$\frac{1}{36} \begin{bmatrix} 2(1+s) & 0 & -2(1+s) & 0 & -2(1-s) & 0 & 2(1-s) & 0 \\ 0 & (1+r) & 0 & (1-r) & 0 & -(1-r) & 0 & -(1+r) \\ 2(1+r) & (1+s) & 2(1-r) & -(1+s) & -2(1-r) & -(1-s) & -2(1+r) & (1-s) \end{bmatrix}$$

Table 6.5

$${}^i_0\mathbf{B}_{NL}$$

= nonlinear strain-displacement transformation matrices

$${}^i_0\mathbf{B}_{NL} = \begin{bmatrix} {}_0h_{1,1} & 0 & {}_0h_{2,1} & 0 & {}_0h_{3,1} & 0 & \cdots & {}_0h_{N,1} & 0 \\ {}_0h_{1,2} & 0 & {}_0h_{2,2} & 0 & {}_0h_{3,2} & 0 & \cdots & {}_0h_{N,2} & 0 \\ 0 & {}_0h_{1,1} & 0 & {}_0h_{2,1} & 0 & {}_0h_{3,1} & \cdots & 0 & {}_0h_{N,1} \\ 0 & {}_0h_{1,2} & 0 & {}_0h_{2,2} & 0 & {}_0h_{3,2} & \cdots & 0 & {}_0h_{N,2} \\ \frac{h_1}{{}_0\bar{x}_1} & 0 & \frac{h_2}{{}_0\bar{x}_1} & 0 & \frac{h_3}{{}_0\bar{x}_1} & 0 & \cdots & \frac{h_N}{{}_0\bar{x}_1} & 0 \end{bmatrix}$$

$${}^i_0\mathbf{B}_{NL} = \frac{1}{6} \begin{bmatrix} (1+s) & 0 & -(1+s) & 0 & -(1-s) & 0 & (1-s) & 0 \\ (1+r) & 0 & (1-r) & 0 & -(1-r) & 0 & -(1+r) & 0 \\ 0 & (1+s) & 0 & -(1+s) & 0 & -(1-s) & 0 & (1-s) \\ 0 & (1+r) & 0 & (1-r) & 0 & -(1-r) & 0 & -(1+r) \end{bmatrix}$$

6.4 Displacement/Pressure Formulation for Large Deformations

(Page 561 – 567)

Why Displacement/Pressure Formulations for Large Deformations?

- A pure displacement-based procedure is not effective for (almost) incompressible analysis
- Material in large deformations (like rubberlike materials or metals under large strain) behave almost incompressible
- Therefore total and updated Lagrangian formulations should be extended to incompressible analysis

Total Lagrangian Formulation pure displacement-based

Equation of motion
(based on the principle of
virtual displacements):

$$\int_{\mathcal{V}_0} {}^{t+\Delta t} \bar{\mathbf{S}}_{ij} \delta {}^{t+\Delta t} \bar{\boldsymbol{\epsilon}}_{ij} d^0 V = {}^{t+\Delta t} \mathcal{R}$$

We define the incremental potential*:

$$d {}^t_0 \bar{W} = {}^t_0 \bar{\mathbf{S}}_{ij} d {}^t_0 \bar{\boldsymbol{\epsilon}}_{ij}$$

Now the principle of virtual
displacement can be rewritten as:

$$\delta \left(\int_{\mathcal{V}_0} {}^t_0 \bar{W} d^0 V \right) = {}^t \mathcal{R}$$

*The overbar denotes that that only
displacement fields where
considered (not pressure fields)

Displacement/Pressure Total Lagrangian Formulation

$$\delta \left(\int_{\mathcal{V}_0} {}^t_0 W \, d^0V \right) = {}^t_0 \mathcal{R}$$

$${}^t_0 W = {}^t_0 \bar{W} + {}^t_0 Q$$

${}^t_0 W$: Total Potential

${}^t_0 \bar{W}$: Potential of the displacements

${}^t_0 Q$: Potential of the pressure

Potential Q

$${}^t_0Q = -\frac{1}{2\kappa} ({}^t\bar{p} - {}^t\tilde{p})^2$$

${}^t\bar{p}$: Pressure of the displacement field

${}^t\tilde{p}$: Total element pressure at time t

κ : bulk modulus

$${}^t u_i = \sum_{k=1}^N h_k {}^t u_i^k; \quad {}^t\tilde{p} = \sum_{i=1}^q g_i {}^t \hat{p}_i$$

$${}^t_0W = {}^t_0\bar{W} - \frac{1}{2\kappa} ({}^t\bar{p} - {}^t\tilde{p})^2$$

Bulk Modulus κ

$$\delta Q = -\frac{1}{2\kappa} (\bar{p} - \tilde{p})^2$$

- The bulk modulus (Kompressionsmodul) is defined as the pressure increase needed to cause a given relative decrease in volume

$$\kappa = -V \frac{\partial P}{\partial V}$$

Linearization

$$\begin{pmatrix} {}^t\mathbf{KUU} & {}^t\mathbf{KUP} \\ {}^t\mathbf{KPU} & {}^t\mathbf{KPP} \end{pmatrix} \begin{pmatrix} \hat{\mathbf{u}} \\ \hat{\mathbf{p}} \end{pmatrix} = \begin{pmatrix} {}^{t+\Delta t}\mathbf{R} \\ \mathbf{0} \end{pmatrix} - \begin{pmatrix} {}^t\mathbf{FU} \\ {}^t\mathbf{FP} \end{pmatrix}$$

Vectors FU, FP contains these Elements

$${}^tFU_i = \frac{\partial}{\partial {}^t\hat{u}_i} \left(\int_{0V} {}^tW d^0V \right) = \int_{0V} {}^tS_{ki} \frac{\partial {}^t\epsilon_{ki}}{\partial {}^t\hat{u}_i} d^0V$$

$${}^tFP_i = \frac{\partial}{\partial {}^t\hat{p}_i} \left(\int_{0V} {}^tW d^0V \right) = \int_{0V} \frac{1}{\kappa} ({}^t\bar{p} - {}^t\tilde{p}) \frac{\partial {}^t\tilde{p}}{\partial {}^t\hat{p}_i} d^0V$$

$$\uparrow$$

$${}^tQ = -\frac{1}{2\kappa} ({}^t\bar{p} - {}^t\tilde{p})^2$$

Linearization

$$\begin{pmatrix} {}^t\mathbf{KUU} & {}^t\mathbf{KUP} \\ {}^t\mathbf{KPU} & {}^t\mathbf{KPP} \end{pmatrix} \begin{pmatrix} \hat{\mathbf{u}} \\ \hat{\mathbf{p}} \end{pmatrix} = \begin{pmatrix} {}^{t+\Delta t}\mathbf{R} \\ \mathbf{0} \end{pmatrix} - \begin{pmatrix} {}^t\mathbf{FU} \\ {}^t\mathbf{FP} \end{pmatrix}$$

Matrices KUU, KUP, KPP, KPU contains these Elements

$${}^tKUU_{ij} = \frac{\partial {}^tFU_i}{\partial {}^t\hat{u}_j} = \int_{0V} {}_0CUU_{kirs} \frac{\partial {}^t_0\epsilon_{ki}}{\partial {}^t\hat{u}_i} \frac{\partial {}^t_0\epsilon_{rs}}{\partial {}^t\hat{u}_j} d^0V + \int_{0V} {}^t_0S_{kl} \frac{\partial^2 {}^t_0\epsilon_{ki}}{\partial {}^t\hat{u}_i \partial {}^t\hat{u}_j} d^0V$$

$${}^tKUP_{ij} = \frac{\partial {}^tFU_i}{\partial {}^t\hat{p}_j} = \frac{\partial {}^tFP_j}{\partial {}^t\hat{u}_i} = {}^tKPU_{ji} = \int_{0V} {}_0CUP_{ki} \frac{\partial {}^t_0\epsilon_{ki}}{\partial {}^t\hat{u}_i} \frac{\partial {}^t\bar{p}}{\partial {}^t\hat{p}_j} d^0V$$

$${}^tKPP_{ij} = \frac{\partial {}^tFP_i}{\partial {}^t\hat{p}_j} = \int_{0V} -\frac{1}{\kappa} \frac{\partial {}^t\bar{p}}{\partial {}^t\hat{p}_i} \frac{\partial {}^t\bar{p}}{\partial {}^t\hat{p}_j} d^0V$$

$${}^i_0 S_{ki} = {}^i_0 \bar{S}_{ki} - \frac{1}{\kappa} ({}^i\bar{p} - {}^i\tilde{p}) \frac{\partial {}^i\bar{p}}{\partial {}^i_0 \epsilon_{ki}}$$

$${}^i_0 CUU_{kirs} = {}^i_0 \bar{C}_{kirs} - \frac{1}{\kappa} \frac{\partial {}^i\bar{p}}{\partial {}^i_0 \epsilon_{ki}} \frac{\partial {}^i\bar{p}}{\partial {}^i_0 \epsilon_{rs}} - \frac{1}{\kappa} ({}^i\bar{p} - {}^i\tilde{p}) \frac{\partial^2 {}^i\bar{p}}{\partial {}^i_0 \epsilon_{ki} \partial {}^i_0 \epsilon_{rs}}$$

$${}^i_0 CUP_{ki} = \frac{1}{\kappa} \frac{\partial {}^i\bar{p}}{\partial {}^i_0 \epsilon_{ki}}$$

$$\frac{\partial {}^i\tilde{p}}{\partial {}^i\hat{p}_i} = g_i$$

$$\frac{\partial {}^i_0 \epsilon_{kl}}{\partial {}^i u_n^L} = \frac{1}{2} ({}^i x_{n,k} {}^i_0 h_{L,l} + {}^i x_{n,l} {}^i_0 h_{L,k})$$

$$\frac{\partial^2 {}^i_0 \epsilon_{kl}}{\partial {}^i u_n^L \partial {}^i u_m^M} = \frac{1}{2} ({}^i_0 h_{L,k} {}^i_0 h_{M,l} + {}^i_0 h_{L,l} {}^i_0 h_{M,k}) \delta_{nm}$$

Displacement/Pressure Updated Lagrangian Formulation

Total Lagrange: $\delta \int_{\mathcal{V}_0} \left({}^i\bar{W} + {}^iQ \right) d^0V = {}^i\mathcal{R}$




Updated Lagrange: $\delta \int_{\mathcal{V}^T} \left({}^t\bar{W} + {}^tQ \right) d^TV = {}^i\mathcal{R}$

$${}^t\bar{W} d^TV = {}^i\bar{W} d^0V$$

$${}^tQ d^TV = {}^iQ d^0V$$

Innovation in UL formulation

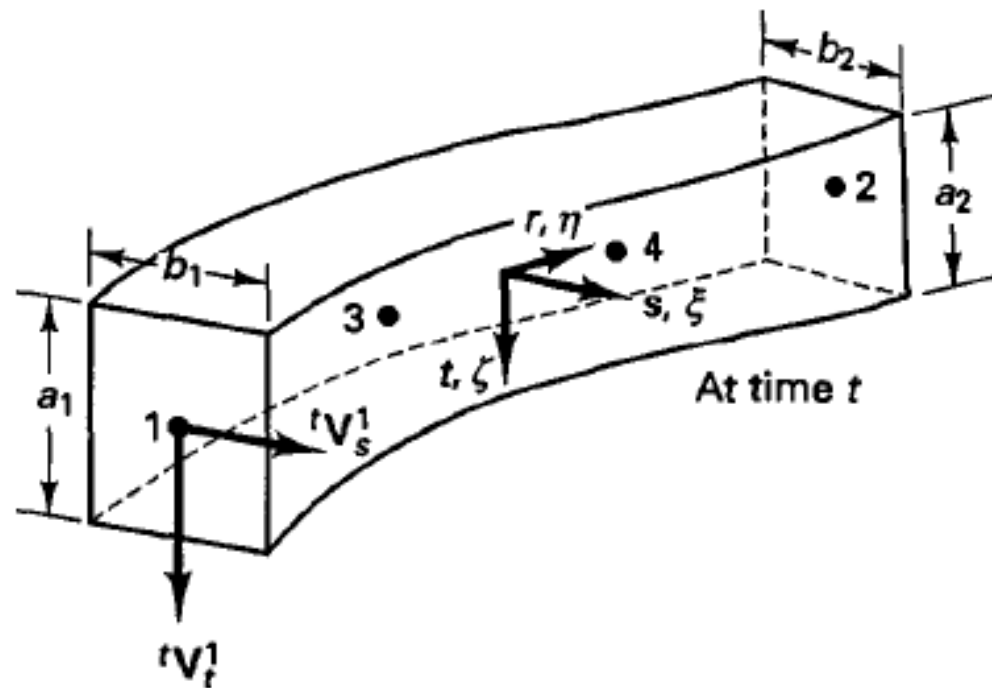
$${}^t\mathbf{Q} d^T V = {}^t_0\mathbf{Q} d^0 V, \quad {}^t_0\mathbf{Q} = -\frac{1}{2\kappa} ({}^t\bar{p} - {}^t\tilde{p})^2$$


$${}^t\mathbf{Q} = -\frac{1}{2\kappa^*} ({}^t\bar{p} - {}^t\tilde{p})^2$$

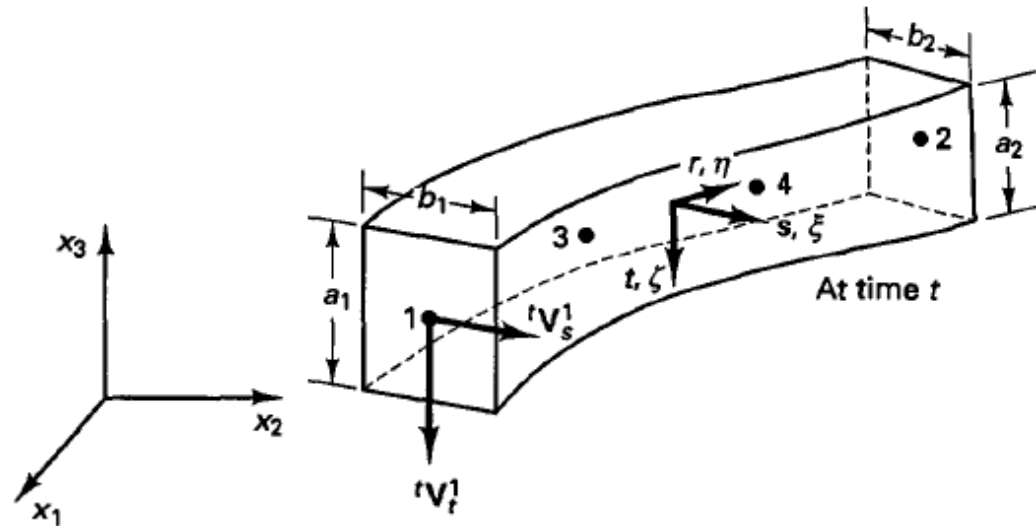
$$\kappa^* = \kappa \frac{d^T V}{d^0 V} = \kappa \det {}^T_0 \mathbf{X}$$

6.5.1 Structural Elements: Beam and Axisymmetric Shell Elements

(Page 568 – 572)



Dimensions of the Beam



$${}^t x_i = \sum_{k=1}^q h_k {}^t x_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k {}^t V_{ti}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k {}^t V_{si}^k$$

The vectors V_s and V_t define the Orientation of the cross-section of the beam: they are normal to the axis of the beam and to each other
The values a and b define the size of the cross-section of the beam

Displacement ${}^t u$ / Incremental Displacement u

$${}^t u_i = {}^t x_i - {}^0 x_i$$

$${}^t x_i = \sum_{k=1}^q h_k {}^t x_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k {}^t V_{ii}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k {}^t V_{si}^k$$

$${}^t u_i = \sum_{k=1}^q h_k {}^t u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k ({}^t V_{ii}^k - {}^0 V_{ii}^k) + \frac{s}{2} \sum_{k=1}^q b_k h_k ({}^t V_{si}^k - {}^0 V_{si}^k)$$

$$u_i = {}^{t+\Delta t} x_i - {}^t x_i$$

$$V_{ii}^k = {}^{t+\Delta t} V_{ii}^k - {}^t V_{ii}^k$$

$$V_{si}^k = {}^{t+\Delta t} V_{si}^k - {}^t V_{si}^k$$

$$u_i = \sum_{k=1}^q h_k u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k V_{ii}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k V_{si}^k$$

Vector V_t and V_s

FEM I:
$$\mathbf{V}_t^k = \boldsymbol{\theta}_k \times {}^0\mathbf{V}_t^k = e_{x1} * (\theta_{x2}^k * {}^0V_{t,x3}^k - \theta_{x3}^k * {}^0V_{t,x2}^k) + e_{x2} \dots$$

$$\mathbf{V}_s^k = \boldsymbol{\theta}_k \times {}^0\mathbf{V}_s^k$$

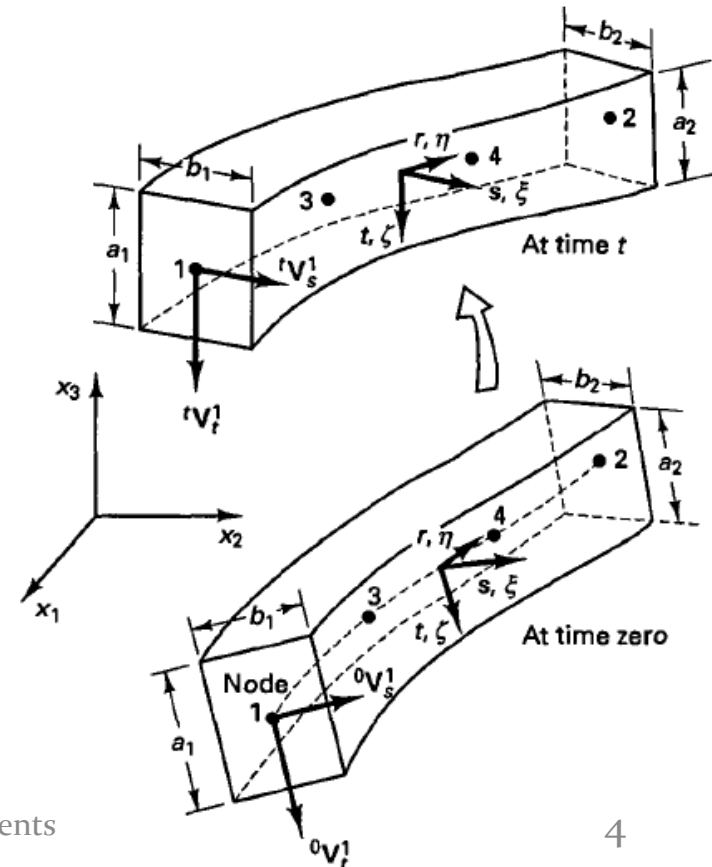
$$\boldsymbol{\theta}_k \times {}^0\mathbf{V}_t^k = \det \begin{bmatrix} e_x & e_y & e_z \\ \theta_x^k & \theta_y^k & \theta_z^k \\ V_{t,x}^k & V_{t,y}^k & V_{t,z}^k \end{bmatrix}$$

$$\boldsymbol{\theta}_k = \begin{bmatrix} \theta_{x1}^k \\ \theta_{x2}^k \\ \theta_{x3}^k \end{bmatrix}$$

$\boldsymbol{\theta}_k$ is a vector listing the nodal point rotations at nodal point k

$$V_{ti}^k = {}^{t+\Delta t}V_{ti}^k - {}^tV_{ti}^k$$

$$V_{si}^k = {}^{t+\Delta t}V_{si}^k - {}^tV_{si}^k$$

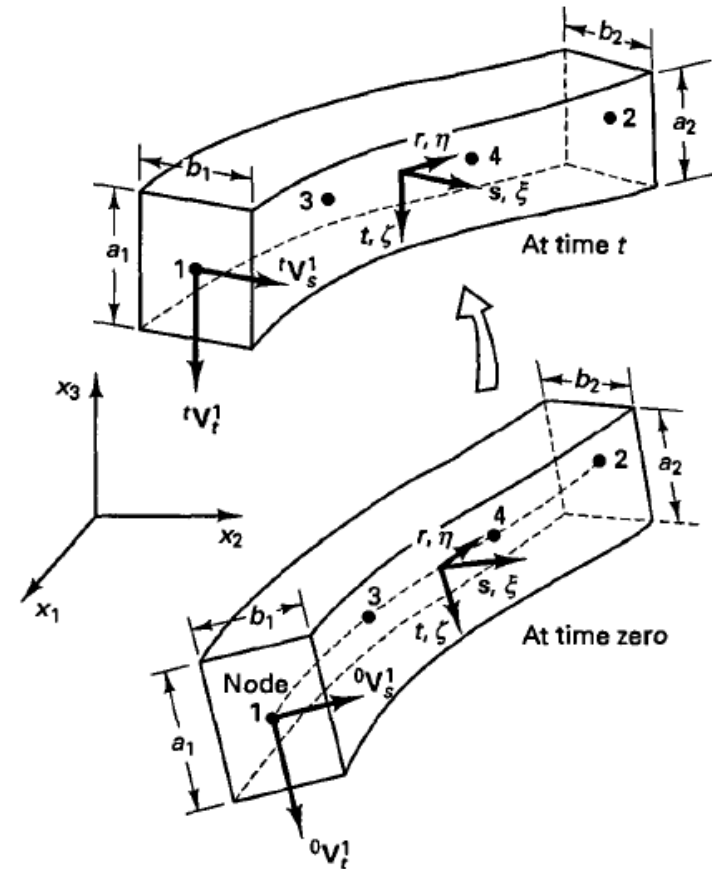


Problem: Large Rotation

-> therefore second-order has to be taken into account
(higher order terms are neglected)

$$\mathbf{V}_t^k = \boldsymbol{\theta}_k \times {}^t\mathbf{V}_t^k + \frac{1}{2} \boldsymbol{\theta}_k \times (\boldsymbol{\theta}_k \times {}^t\mathbf{V}_t^k)$$

$$\mathbf{V}_s^k = \boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k + \frac{1}{2} \boldsymbol{\theta}_k \times (\boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k)$$



Result

Now the Equation can be solved

$${}^t\mathbf{K} \begin{bmatrix} \vdots \\ \mathbf{u}_k \\ \boldsymbol{\theta}_k \\ \vdots \end{bmatrix} = {}^{t+\Delta t}\mathbf{R} - {}^t\mathbf{F}$$

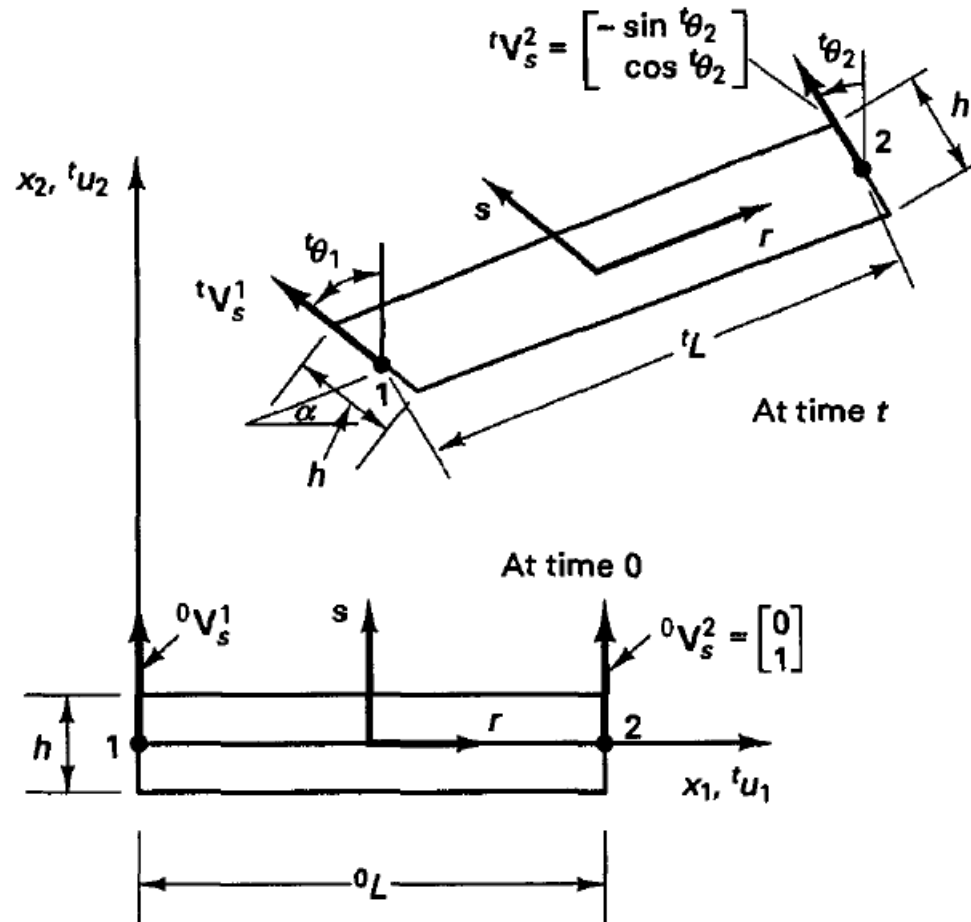
$${}^{t+\Delta t}\mathbf{u}_k = {}^t\mathbf{u}_k + \mathbf{u}_k$$

$${}^{t+\Delta t}\mathbf{V}_i^k = {}^t\mathbf{V}_i^k + \int_{\boldsymbol{\theta}_k} d\boldsymbol{\theta}_k \times {}^t\mathbf{V}_i^k$$

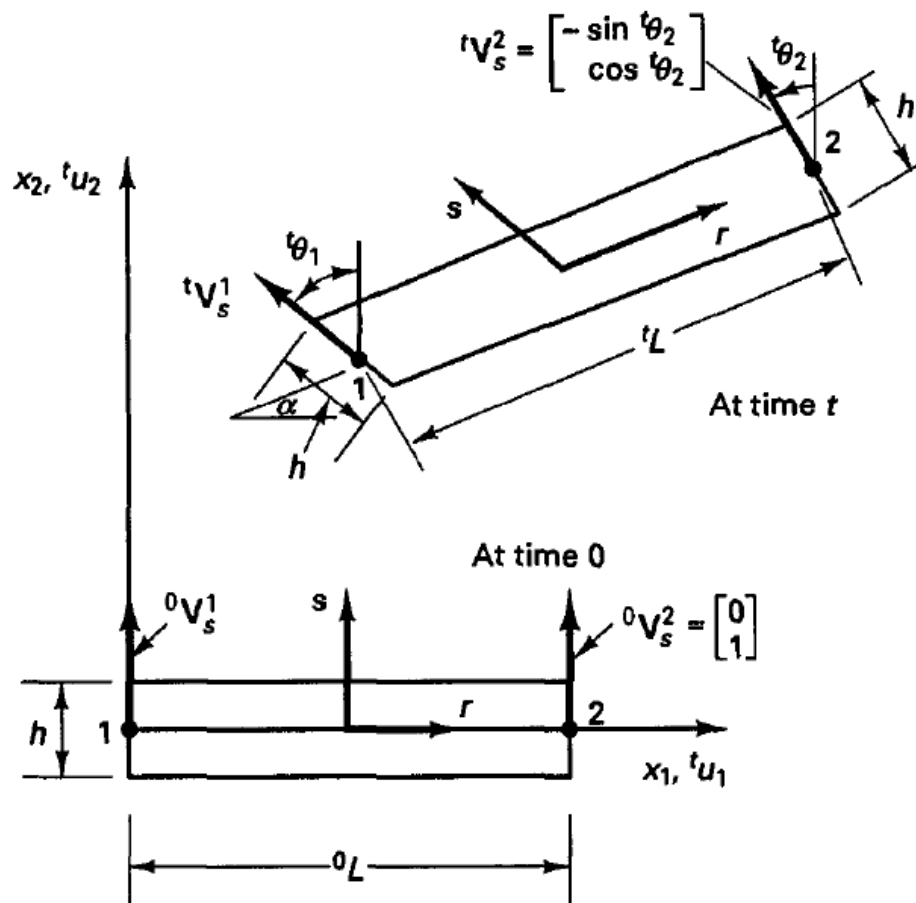
$${}^{t+\Delta t}\mathbf{V}_s^k = {}^t\mathbf{V}_s^k + \int_{\boldsymbol{\theta}_k} d\boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k$$

Example 6.20

Two Node Beam Element in Large Displacement and Rotation



Evaluate the coordinate and displacement interpolations and derivatives that are required for the calculation of the strain-displacement matrices of the UL and TL formulations

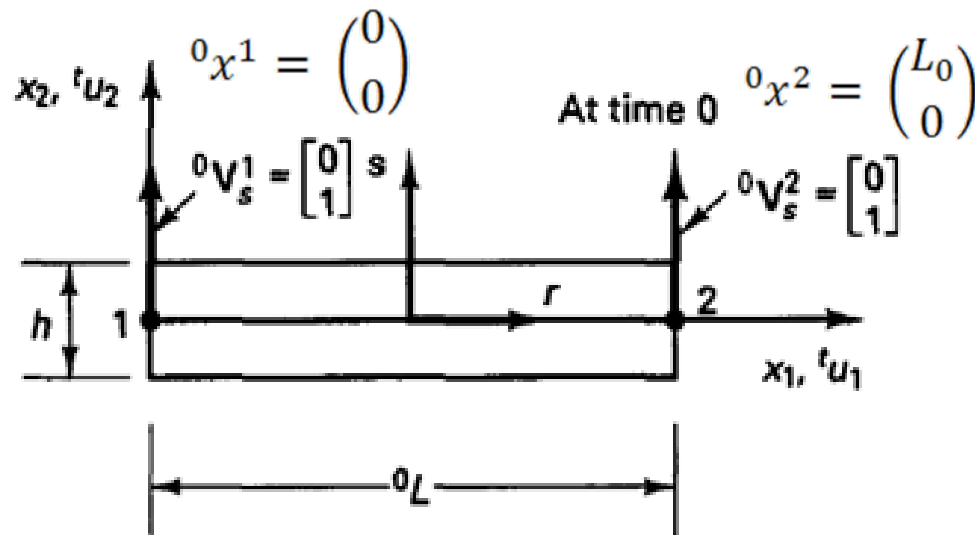


First we have to calculate:

$${}^t x_1 \quad {}^t x_2$$

$${}^0 x_1 \quad {}^0 x_2$$

$${}^t x_i = \sum_{k=1}^q h_k {}^t x_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k {}^t V_{ti}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k {}^t V_{si}^k$$



Two-Node Element:

$$h_1 = \frac{1}{2}(1 - r)$$

$$h_2 = \frac{1}{2}(1 + r)$$

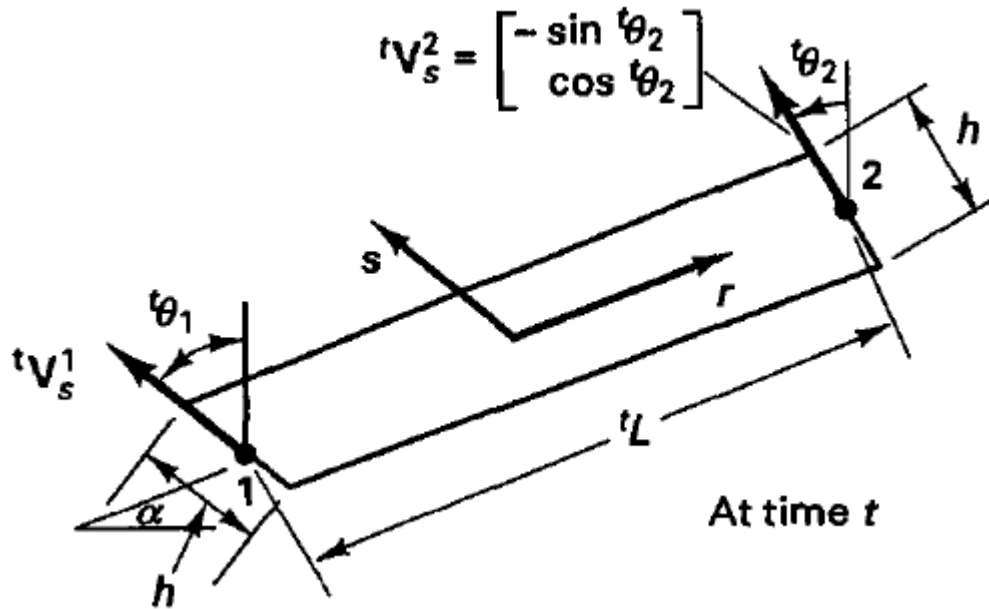
$${}^0 x_1 = h_1 {}^0 x_1^1 + h_2 {}^0 x_1^2 + \frac{s}{2} b_1 h_1 {}^0 V_{s1}^1 + \frac{s}{2} b_2 h_2 {}^0 V_{s1}^2$$

$${}^0 x_1 = \left(\frac{1+r}{2} \right) {}^0 L$$

$${}^0 x_2 = \frac{sh}{2}$$

Coordinates

$${}^t x_i = \sum_{k=1}^q h_k {}^t x_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k {}^t V_{ti}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k {}^t V_{si}^k$$

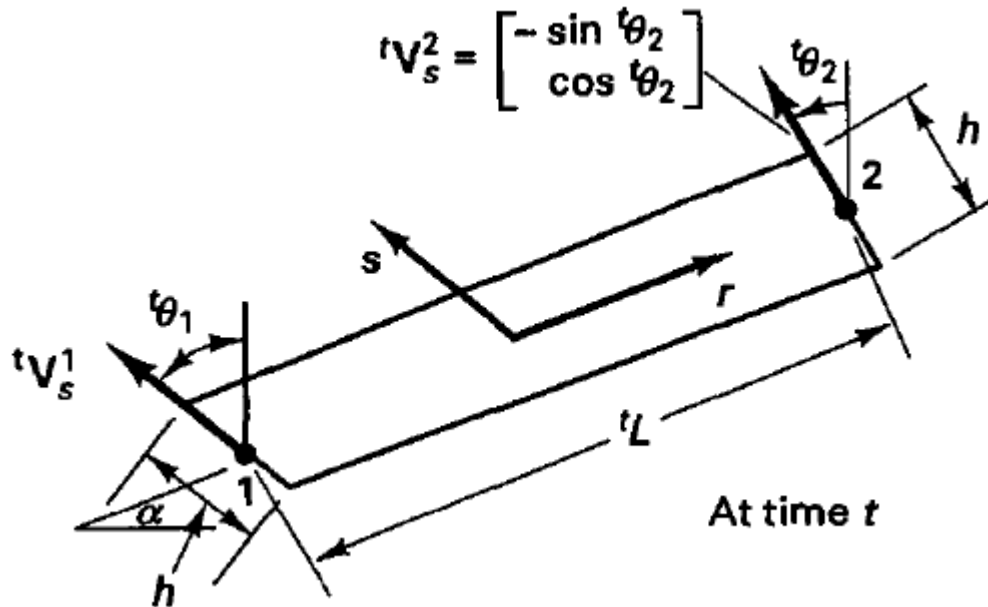


$${}^t x_1 = \left(\frac{1-r}{2}\right) {}^t x_1^1 + \left(\frac{1+r}{2}\right) {}^t x_1^2 - \frac{sh}{2} \left(\frac{1-r}{2}\right) \sin {}^t \theta_1 - \frac{sh}{2} \left(\frac{1+r}{2}\right) \sin {}^t \theta_2$$

$${}^t x_2 = \left(\frac{1-r}{2}\right) {}^t x_2^1 + \left(\frac{1+r}{2}\right) {}^t x_2^2 + \frac{sh}{2} \left(\frac{1-r}{2}\right) \cos {}^t \theta_1 + \frac{sh}{2} \left(\frac{1+r}{2}\right) \cos {}^t \theta_2$$

Displacement

$${}^t u_i = \sum_{k=1}^q h_k {}^t u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k ({}^t V_{ii}^k - {}^0 V_{ii}^k) + \frac{s}{2} \sum_{k=1}^q b_k h_k ({}^t V_{si}^k - {}^0 V_{si}^k)$$



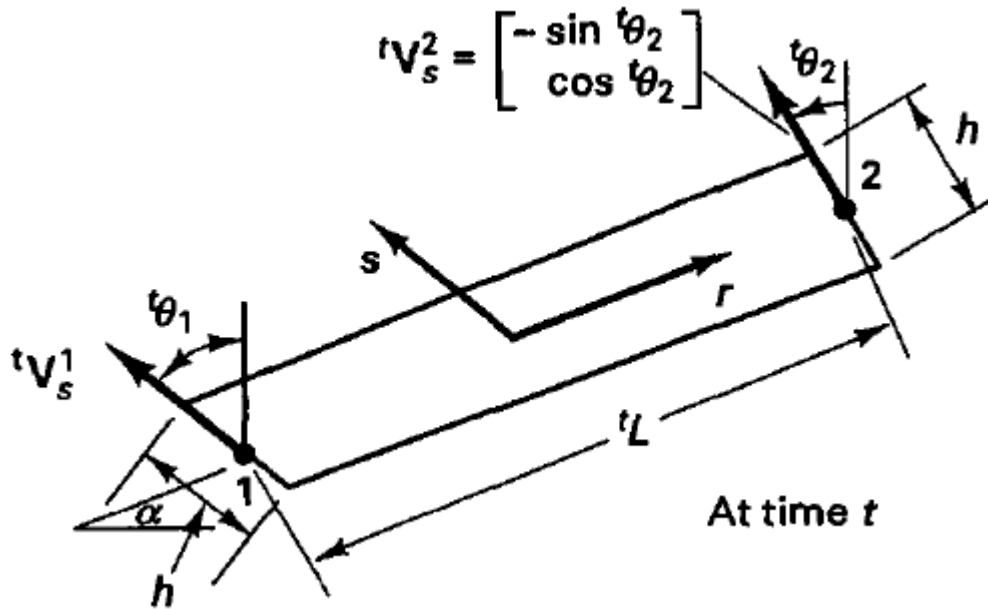
$${}^t u_i = {}^t x_i - {}^0 x_i$$

$${}^t u_1 = \left(\frac{{}^t x_1^1 + {}^t x_1^2 - {}^0 L}{2} \right) + \left(\frac{{}^t x_1^2 - {}^t x_1^1 - {}^0 L}{2} \right) r - \frac{sh}{2} \left[\left(\frac{1-r}{2} \right) \sin \theta_1 + \left(\frac{1+r}{2} \right) \sin \theta_2 \right]$$

$${}^t u_2 = \left(\frac{{}^t x_2^1 + {}^t x_2^2}{2} \right) + \left(\frac{{}^t x_2^2 - {}^t x_2^1}{2} \right) r + \frac{sh}{2} \left[\left(\frac{1-r}{2} \right) \cos \theta_1 + \left(\frac{1+r}{2} \right) \cos \theta_2 - 1 \right]$$

Incremental Displacement

$$u_i = \sum_{k=1}^q h_k u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k V_{ti}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k V_{si}^k$$

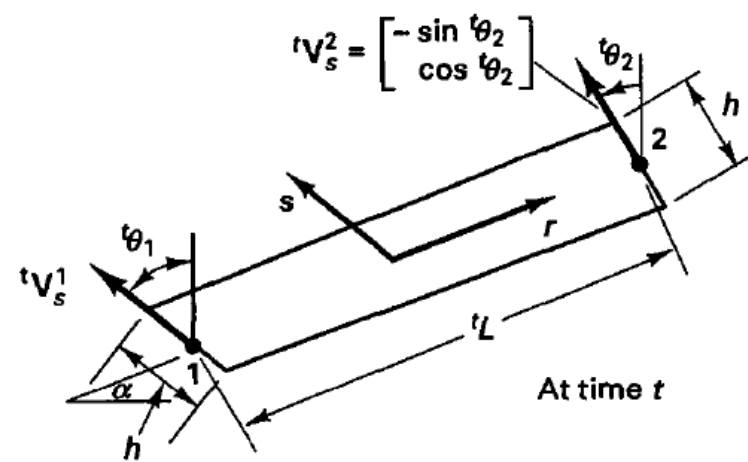


second-order terms

$$V_i^k = \theta_k \times {}^tV_i^k + \frac{1}{2} \theta_k \times (\theta_k \times {}^tV_i^k)$$

$$V_s^k = \theta_k \times {}^tV_s^k + \frac{1}{2} \theta_k \times (\theta_k \times {}^tV_s^k)$$

$$\mathbf{V}_s^k = \boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k + \frac{1}{2} \boldsymbol{\theta}_k \times (\boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k)$$



$$\boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k = \det \begin{bmatrix} e_{x1} & e_{x2} & e_{x3} \\ 0 & 0 & \theta_k \\ -\sin({}^t\theta_k) & \cos({}^t\theta_k) & 0 \end{bmatrix} = \begin{pmatrix} -\theta_k \cos({}^t\theta_k) \\ -\theta_k \sin({}^t\theta_k) \\ 0 \end{pmatrix}$$

$$\boldsymbol{\theta}_k \times (\boldsymbol{\theta}_k \times {}^t\mathbf{V}_s^k) = \det \begin{bmatrix} e_{x1} & e_{x2} & e_{x3} \\ 0 & 0 & \theta_k \\ -\theta_k \cos({}^t\theta_k) & -\theta_k \sin({}^t\theta_k) & 0 \end{bmatrix} = \begin{pmatrix} \theta_k^2 \sin({}^t\theta_k) \\ -\theta_k^2 \cos({}^t\theta_k) \\ 0 \end{pmatrix}$$

$$\mathbf{V}_s^k = \begin{pmatrix} -\theta_k \cos({}^t\theta_k) + \frac{1}{2} \theta_k^2 \sin({}^t\theta_k) \\ -\theta_k \sin({}^t\theta_k) - \frac{1}{2} \theta_k^2 \cos({}^t\theta_k) \\ 0 \end{pmatrix}$$

$$u_i = \sum_{k=1}^q h_k u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k V_{ri}^k + \frac{s}{2} \sum_{k=1}^q b_k h_k V_{si}^k$$

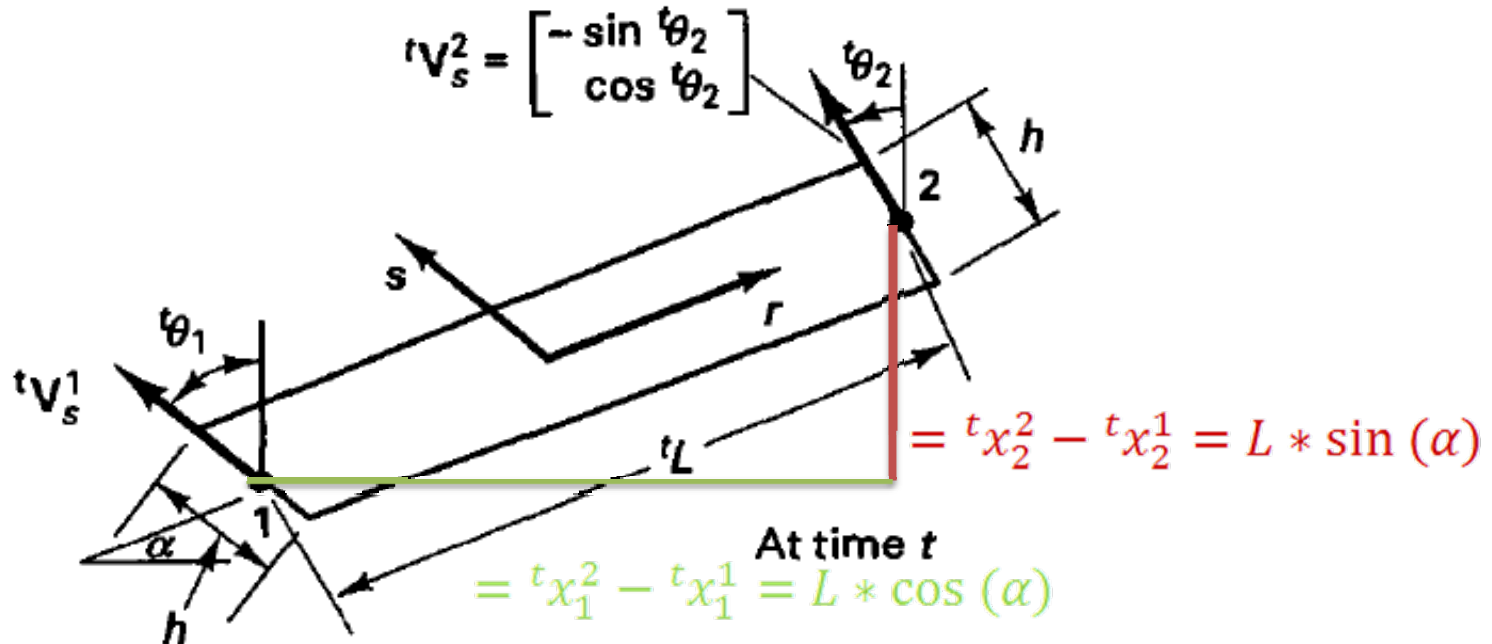
$$V_s^2 = \begin{pmatrix} -\theta_k \cos({}^t\theta_k) + \frac{1}{2} \theta_k^2 \sin({}^t\theta_k) \\ -\theta_k \sin({}^t\theta_k) - \frac{1}{2} \theta_k^2 \cos({}^t\theta_k) \\ 0 \end{pmatrix}$$

$$u_1 = \frac{1-r}{2} u_1^1 + \frac{1+r}{2} u_1^2 + \frac{sh}{2} \left(\frac{1-r}{2} \right) \left[(-\cos {}^t\theta_1) \theta_1 + \frac{1}{2} \sin {}^t\theta_1 (\theta_1)^2 \right] \\ + \frac{sh}{2} \left(\frac{1+r}{2} \right) \left[(-\cos {}^t\theta_2) \theta_2 + \frac{1}{2} \sin {}^t\theta_2 (\theta_2)^2 \right]$$

$$\begin{aligned}
 u_2 = & \frac{1-r}{2} u_2^1 + \frac{1+r}{2} u_2^2 + \frac{sh}{2} \left(\frac{1-r}{2} \right) \left[(-\sin \theta_1) \theta_1 - \frac{1}{2} \cos \theta_1 (\theta_1)^2 \right] \\
 & + \frac{sh}{2} \left(\frac{1+r}{2} \right) \left[(-\sin \theta_2) \theta_2 - \frac{1}{2} \cos \theta_2 (\theta_2)^2 \right]
 \end{aligned}$$

Derivatives

$${}^t x_1 = \left(\frac{1-r}{2}\right) {}^t x_1^1 + \left(\frac{1+r}{2}\right) {}^t x_1^2 - \frac{sh}{2} \left(\frac{1-r}{2}\right) \sin {}^t \theta_1 - \frac{sh}{2} \left(\frac{1+r}{2}\right) \sin {}^t \theta_2$$



$$\frac{\partial {}^t x_1}{\partial r} = \frac{L \cos \alpha}{2} - \frac{sh}{4} (\sin {}^t \theta_2 - \sin {}^t \theta_1)$$

Derivatives

$${}'x_1 = \left(\frac{1-r}{2}\right) {}'x_1^1 + \left(\frac{1+r}{2}\right) {}'x_1^2 - \frac{sh}{2}\left(\frac{1-r}{2}\right) \sin {}'\theta_1 - \frac{sh}{2}\left(\frac{1+r}{2}\right) \sin {}'\theta_2$$

$${}'x_2 = \left(\frac{1-r}{2}\right) {}'x_2^1 + \left(\frac{1+r}{2}\right) {}'x_2^2 + \frac{sh}{2}\left(\frac{1-r}{2}\right) \cos {}'\theta_1 + \frac{sh}{2}\left(\frac{1+r}{2}\right) \cos {}'\theta_2$$

$$\frac{\partial {}'x_1}{\partial r} = \frac{L \cos \alpha}{2} - \frac{sh}{4}(\sin {}'\theta_2 - \sin {}'\theta_1)$$

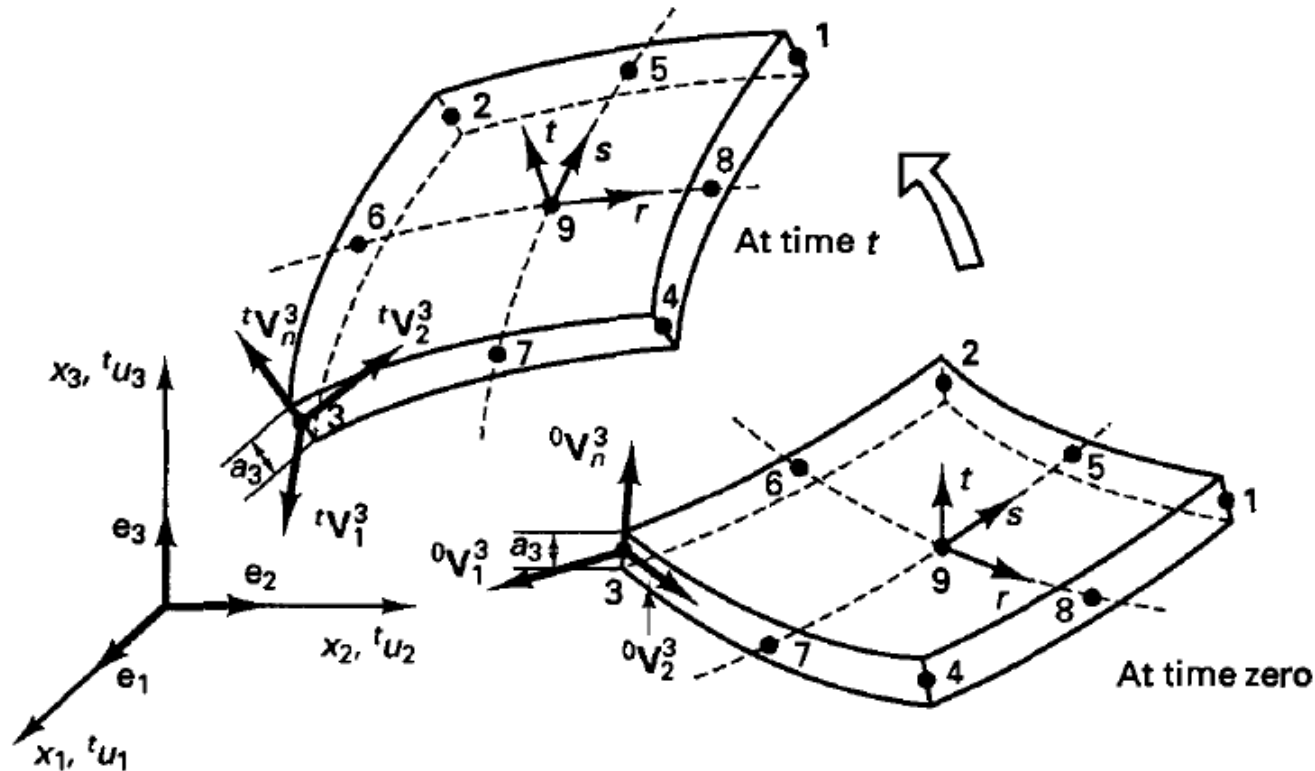
$$\frac{\partial {}'x_1}{\partial s} = \left(-\frac{h}{2}\right) \left[\left(\frac{1-r}{2}\right) \sin {}'\theta_1 + \left(\frac{1+r}{2}\right) \sin {}'\theta_2 \right]$$

$$\frac{\partial {}'x_2}{\partial r} = \frac{L \sin \alpha}{2} + \frac{sh}{4}(\cos {}'\theta_2 - \cos {}'\theta_1)$$

$$\frac{\partial {}'x_2}{\partial s} = \frac{h}{2} \left[\left(\frac{1-r}{2}\right) \cos {}'\theta_1 + \left(\frac{1+r}{2}\right) \cos {}'\theta_2 \right]$$

6.5.2 Structural Elements: Structural Elements: Plate and General Shell Elements

(Page 575 – 578)



We now extend the theory for beams for general shell elements

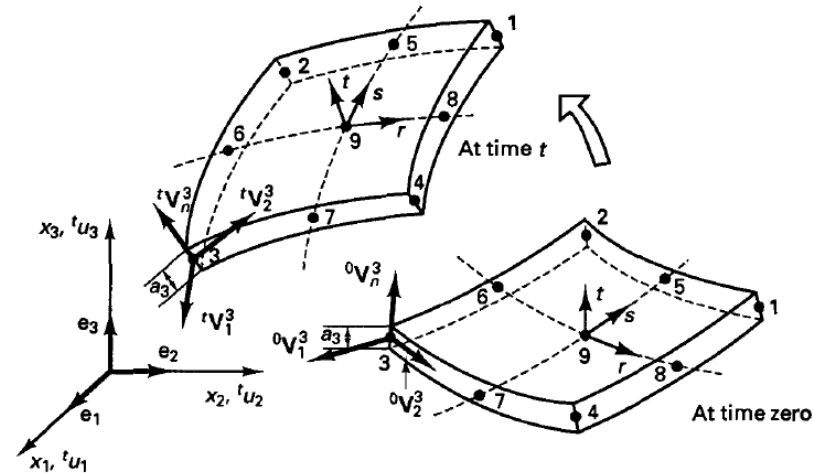
$${}^t x_i = \sum_{k=1}^q h_k {}^t x_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k {}^t V_{ni}^k$$

$$\downarrow$$

$${}^t u_i = {}^t x_i - {}^0 x_i$$

$$\downarrow$$

$${}^t u_i = \sum_{k=1}^q h_k {}^t u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k ({}^t V_{ni}^k - {}^0 V_{ni}^k)$$



$$u_i = {}^{t+\Delta t} x_i - {}^t x_i \quad V_{ni}^k = {}^{t+\Delta t} V_{ni}^k - {}^t V_{ni}^k$$

$$\downarrow \quad \downarrow$$

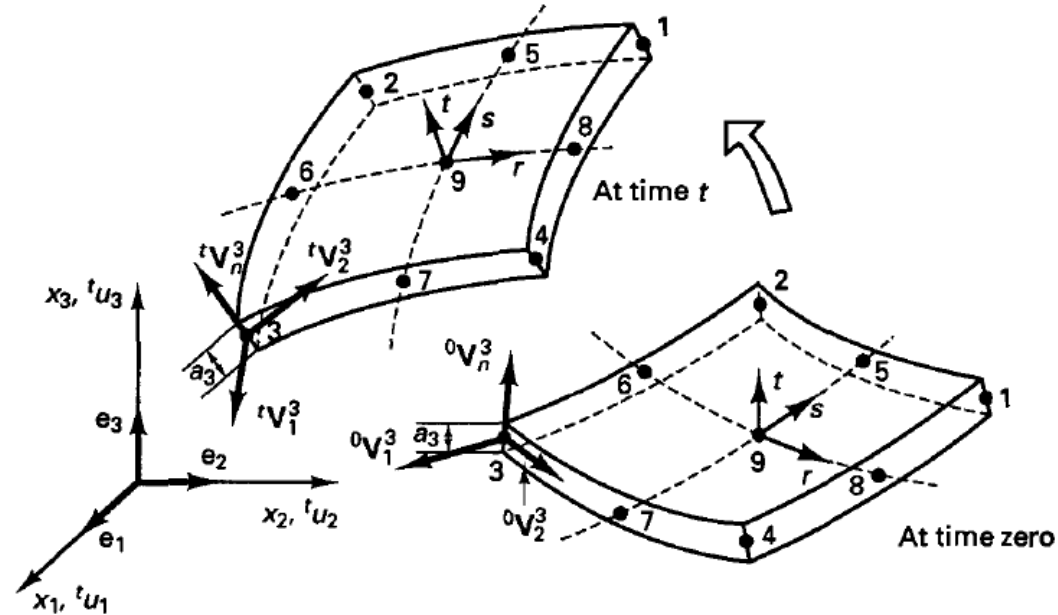
$$u_i = \sum_{k=1}^q h_k u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k V_{ni}^k$$

Rotation

$${}^0\mathbf{V}_1^k = \frac{\mathbf{e}_2 \times {}^0\mathbf{V}_n^k}{\|\mathbf{e}_2 \times {}^0\mathbf{V}_n^k\|_2}$$

$${}^0\mathbf{V}_2^k = {}^0\mathbf{V}_n^k \times {}^0\mathbf{V}_1^k$$

${}^i\mathbf{V}_1^k$ and ${}^i\mathbf{V}_2^k$ are defined at time 0 Normal Vector



α and β are the Rotation of the Vector ${}^i\mathbf{V}_n^k$ about ${}^i\mathbf{V}_1^k$ and ${}^i\mathbf{V}_2^k$

$$\mathbf{V}_n^k = -{}^i\mathbf{V}_2^k \alpha_k + {}^i\mathbf{V}_1^k \beta_k - \frac{1}{2}(\alpha_k^2 + \beta_k^2) {}^i\mathbf{V}_n^k$$

$${}^{t+\Delta t}\mathbf{V}_n^k = {}^t\mathbf{V}_n^k + \int_{\alpha_k, \beta_k} -{}^t\mathbf{V}_2^k d\alpha_k + {}^t\mathbf{V}_1^k d\beta_k$$

$$u_i = \sum_{k=1}^q h_k u_i^k + \frac{t}{2} \sum_{k=1}^q a_k h_k \left[-{}^t\mathbf{V}_{2i}^k \alpha_k + {}^t\mathbf{V}_{1i}^k \beta_k - \frac{1}{2} (\alpha_k^2 + \beta_k^2) {}^t\mathbf{V}_{ni}^k \right]$$

Note: the procedure for computing the displacements, etc. is now the same as always.

But as seen in FEM I there can be problems with shear-locking. Therefore Mixed Interpolation should be considered.