

The Finite Element Method for the Analysis of Non-Linear and Dynamic Systems



Prof. Dr. Michael Havbro Faber

Swiss Federal Institute of Technology ETH Zurich, Switzerland Page 1



Contents of Today's Lecture

- Motivation, overview and organization of the course
- Introduction to non-linear analysis
- Formulation of the continuum mechanics incremental equations of motion



Motivation

In FEM 1 we learned about the steady state analysis of linear systems

however,

the systems we are dealing with in structural engineering are generally not steady state and also not linear

We must be able to assess the need for a particular type of analysis and we must be able to perform it





Motivation

What kind of problems are not steady state and linear?

E.g. when the:

material behaves non-linearly

deformations become big (p- Δ effects)

loads vary fast compared to the eigenfrequencies of the structure

General feature: Response becomes load path dependent



Motivation

What is the "added value" of being able to assess the non-linear non-steady state response of structures ?

- **E.g.** assessing the:
- structural response of structures to extreme events (rock-fall, earthquake, hurricanes)
- performance (failures and deformations) of soils
- verifying simple models



• Collapse Analysis of the World Trade Center





• Collapse Analysis of the World Trade Center





• Analysis of ultimate collapse capacity of jacket structure







Analysis of ultimate collapse capacity of jacket structure





Analysis of soil performance





• Analysis of bridge response





Steady state problems (Linear/Non-linear):

The response of the system does not change over time

Propagation problems (Linear/Non-linear):

The response of the system changes over time $M\ddot{U}(t) + C\dot{U}(t) + KU(t) = R(t)$

Eigenvalue problems:

No unique solution to the response of the system $\mathbf{A}\mathbf{v}=\boldsymbol{\lambda}\mathbf{B}\mathbf{v}$

Method of Finite Elements II

 $\mathbf{K}\mathbf{U} = \mathbf{R}$



Organization

The lectures will be given by:

M. H. Faber

Exercises will be organized/attended by:

Jianjun Qin

By appointment, HIL E13.1.





• Organization

PowerPoint files with the presentations will be uploaded on our homepage one day in advance of the lectures

http://www.ibk.ethz.ch/fa/education/FE_II

The lecture as such will follow the book:

"Finite Element Procedures" by K.J. Bathe, Prentice Hall, 1996



• Overview

Date	Pages	Subject
18.09.2009	485-502	 Non-linear Finite Element Calculations in solids and structural mechanics Introduction to non-linear calculations
		 The incremental approach to continuum mechanics
25.09.2009	502-528	Non-linear Finite Element Calculations in solids and structural mechanics
		 Deformation gradients, strain and stress tensors
		 The Langrangian formulation – only material non-linearity
02.10.2009	538-548	Non-linear Finite Element Calculations in solids and structural mechanics
		- Displacement based iso-parametric finite elements in continuum mechanics
09.10.2009	548-560	Non-linear Finite Element Calculations in solids and structural mechanics - Displacement based iso-parametric finite elements in continuum mechanics



• Overview

Non-linear Finite Element Calculations in solids and structural mechanics 16.10.2009 561-578 Total Langrangian formulation Extended Langrangian formulation Structural elements Non-linear Finite Element Calculations in solids and structural mechanics 23.10.2009 581-617 Introduction to constitutive relations Non-linear constitutive relations Non-linear Finite Element Calculations in solids and structural mechanics 30.10.2009 622-640 Contact problems -Practical considerations Dynamical Finite Element Calculations 06.11.2009 768-784 Introduction -Direct integration methods



• Overview

13.11.2009	785-800	Dynamical Finite Element Calculations - Mode superposition
20.11.2009	801-815	Dynamical Finite Element Calculations - Analysis of direct integration methods
27.11.2009	824-830	Dynamical Finite Element Calculations - Solution of dynamical non-linear problems
04.12.2009	887-910	Solution of Eigen value problems - The vector iteration method
11.12.2009	911-937	Solution of Eigen value problems - The transformation method
18.12.2009		Introduction to FEM-software
Method of Finite Elem	ents II	



• Previously we considered the solution of the following linear and static problem:

 $\mathbf{K}\mathbf{U} = \mathbf{R}$

for these problems we have the convenient property of linearity, i.e:

$$\mathbf{K}\mathbf{U} = \lambda \mathbf{R}, \quad \lambda = 1$$
$$\bigcup$$
$$\mathbf{U}^* = \lambda \mathbf{U}, \quad \lambda \neq 1$$

If this is not the case we are dealing with a non-linear problem!



• Previously we considered the solution of the following linear and static problem:

 $\mathbf{K}\mathbf{U} = \mathbf{R}$

we assumed:

small displacements when developing the stiffness matrix K and the load vector R, because we performed all integrations over the original element volume

that the B matrix is constant independent of element displacements

the stress-strain matrix C is constant

boundary constraints are constant



• Classification of non-linear analyses

Type of analysis	Description	Typical	Stress and strain	
		formulation used	measures used	
Materially-nonlinear	Infinitesimal	Materially-	Engineering strain	
only	displacements and	nonlinear-only	and stress	
	strains; stress train	(MNO)		
	relation is non-			
	linear			
Large	Displacements and	Total Lagrange (TL)	Second Piola-	
displacements, large	rotations of fibers		Kirchoff stress,	
rotations but small	are large; but fiber		Green-Lagrange	
strains	extensions and		strain	
	angle changes			
	between fibers are	Updated Lagrange	Cauchy stress,	
	small; stress strain	(UL)	Almansi strain	
	relationship may be			
	linear or non-linear			
Large	Displacements and	Total Lagrange (TL)	Second Piola-	
displacements, large	rotations of fibers		Kirchoff stress,	
rotations and large	are large; fiber		Green-Lagrange	
strains	extensions and		strain	
	angle changes	Updated Lagrange		
	between fibers may	(UL)	Cauchy stress,	
	also be large; stress		Logarithmic strain	
	strain relationship			
	may be linear or			
	non-linear			



• Classification of non-linear analyses





• Classification of non-linear analyses





• Classification of non-linear analyses



Large displacements and large rotations but small strains (linear or nonlinear material behavior)



• Classification of non-linear analyses





• Classification of non-linear analyses



Chang in boundary conditions













• Example: Simple bar structure





$${}^{t}\varepsilon_{a} = \frac{{}^{t}u}{L_{a}}, {}^{t}\varepsilon_{b} = -\frac{{}^{t}u}{L_{b}}$$
$${}^{t}R + {}^{t}\sigma_{b}A = {}^{t}\sigma_{a}A$$
$${}^{t}\varepsilon = \frac{{}^{t}\sigma}{E} \text{ (elastic region)}$$
$${}^{t}\varepsilon = \varepsilon_{Y} + \frac{{}^{t}\sigma - \sigma_{Y}}{E_{T}} \text{ (plastic region)}$$
$$\Delta\varepsilon = \frac{\Delta\sigma}{E} \text{ (unloading)}$$

Both sections elastic

$${}^{t}R = EA^{t}u(\frac{1}{L_{a}} + \frac{1}{L_{b}}) \Longrightarrow {}^{t}u = \frac{{}^{t}R}{3 \cdot 10^{6}}$$
$$\sigma_{a} = \frac{{}^{t}R}{3A}, \sigma_{b} = -\frac{2}{3}\frac{{}^{t}R}{A}$$

a11

Introduction to non-linear analysis



Section a is elastic while section b is plastic





• What did we learn from the example?

The basic problem in general nonlinear analysis is to find a state of equilibrium between externally applied loads and element nodal forces

$${}^{t}\mathbf{R} - {}^{t}\mathbf{F} = 0$$

$${}^{t}\mathbf{R} - {}^{t}\mathbf{F} = 0$$

$${}^{t}\mathbf{R} = {}^{t}\mathbf{R}_{B} + {}^{t}\mathbf{R}_{S} + {}^{t}\mathbf{R}_{C}$$

$${}^{t}\mathbf{F} = {}^{t}\mathbf{R}_{I}$$
$${}^{t}\mathbf{F} = \sum_{m} \int_{{}^{t}V^{(m)}} {}^{t}\mathbf{B}^{(m)T} {}^{t}\tau^{(m)} {}^{t}dV^{(m)}$$

includes implicitly also dynamic analysis!



The basic approach in incremental analysis is

$$t+\Delta t \mathbf{R} - t+\Delta t \mathbf{F} = 0$$

assuming that ${}^{t+\Delta t}\mathbf{R}$ is independent of the deformations we have ${}^{t+\Delta t}\mathbf{F} = {}^{t}\mathbf{F} + \mathbf{F}$

We know the solution ^tF at time t and F is the increment in the nodal point forces corresponding to an increment in the displacements and stresses from time t to time $t+\Delta t$ this we can approximate by





The basic approach in incremental analysis is

We may now substitute the tangent stiffness matrix into the equibrium relation

$${}^{t}\mathbf{K}\mathbf{U} = {}^{t+\Delta t}\mathbf{R} - {}^{t}\mathbf{F}$$

$$\Downarrow$$

$${}^{t+\Delta t}\mathbf{U} = {}^{t}\mathbf{U} + \mathbf{U}$$

which gives us a scheme for the calculation of the displacements

the exact displacements at time $t+\Delta t$ correspond to the applied loads at $t+\Delta t$ however we only determined these approximately as we used a tangent stiffness matrix – thus we may have to iterate to find the solution



The basic approach in incremental analysis is

We may use the Newton-Raphson iteration scheme to find the equilibrium within each load increment

 $^{t+\Delta t}\mathbf{K}^{(i-1)}\Delta \mathbf{U}^{(i)} = {}^{t+\Delta t}\mathbf{R} - {}^{t+\Delta t}\mathbf{F}^{(i-1)}$ (out of balance load vector)

$$^{t+\Delta t}\mathbf{U}^{(i)} = {}^{t+\Delta t}\mathbf{U}^{(i-1)} + \Delta\mathbf{U}^{(i)}$$

with initial conditions

 $^{t+\Delta t}\mathbf{U}^{(0)} = {}^{t}\mathbf{U}; \quad {}^{t+\Delta t}\mathbf{K}^{(0)} = {}^{t}\mathbf{K}; \quad {}^{t+\Delta t}\mathbf{F}^{(0)} = {}^{t}\mathbf{F}$





• The basic approach in incremental analysis is

It may be expensive to calculate the tangent stiffness matrix and,

in the Modified Newton-Raphson iteration scheme it is thus only calculated in the beginning of each new load step

in the quasi-Newton iteration schemes the secant stiffness matrix is used instead of the tangent matrix



• We look at the example again – simple bar (two load steps)

$$({}^{t}K_{a} + {}^{t}K_{b})\Delta u^{(i)} = {}^{t+\Delta t}R - ({}^{t+\Delta t}F_{a}^{(i-1)} - {}^{t+\Delta t}F_{b}^{(i-1)})$$

$${}^{t+\Delta t}u^{(i)} = {}^{t+\Delta t}u^{(i-1)} + \Delta u^{(i)}$$

with initial conditions

$${}^{t+\Delta t}u^{(0)} = {}^{t}u; {}^{t+\Delta t}F_{a}^{(0)} = {}^{t}F_{a} {}^{t+\Delta t}F_{b}^{(0)} = {}^{t}F_{b}$$
$${}^{t}K_{a} = \frac{{}^{t}CA}{L_{a}}; {}^{t}K_{b} = \frac{{}^{t}CA}{L_{b}}$$
$${}^{t}C \begin{cases} = E & \text{if section is elastic} \\ = E_{T} & \text{if section is plastic} \end{cases}$$



• We look at the example again – simple bar

Load step 1:
$$t = 1$$
:
 ${}^{0}K_{a} + {}^{0}K_{b}\Delta u^{(1)} = {}^{1}R - {}^{1}F_{a}^{(0)} - {}^{1}F_{b}^{(0)}$
 \downarrow
 $\Delta u^{(1)} = \frac{2 \times 10^{4}}{10^{7}(\frac{1}{10} + \frac{1}{5})} = 6.6667 \times 10^{-3}$
Iteration 1: $(i = 1)$
 ${}^{1}u^{(1)} = {}^{1}u^{(0)} + \Delta u^{(1)} = 6.6667 \times 10^{-3}$
 ${}^{1}\varepsilon_{a}^{(1)} = \frac{{}^{1}u^{(1)}}{L_{a}} = 6.6667 \times 10^{-4} < \varepsilon_{Y}$ (elastic section!)
 ${}^{1}\varepsilon_{b}^{(1)} = \frac{{}^{1}u^{(1)}}{L_{b}} = 1.3333 \times 10^{-3} < \varepsilon_{Y}$ (elastic section!)
 ${}^{1}F_{a}^{(1)} = 6.6667 \times 10^{3};$ ${}^{1}F_{b}^{(1)} = 1.3333 \times 10^{4}$
 $({}^{0}K_{a} + {}^{0}K_{b})\Delta u^{(2)} = {}^{1}R - {}^{1}F_{a}^{(1)} - {}^{1}F_{b}^{(1)} = 0$

Convergence in one iteration!

 $^{1}u = 6.6667 \times 10^{-3}$



• We look at the example again – simple bar

Load step 2:
$$t = 2$$
:
 $({}^{1}K_{a} + {}^{1}K_{b})\Delta u^{(1)} = {}^{2}R - {}^{2}F_{a}^{(0)} - {}^{2}F_{b}^{(0)}$
 \downarrow
 $\Delta u^{(1)} = \frac{(4 \times 10^{4}) - (6.6667 \times 10^{3}) - (1.333 \times 10^{4})}{10^{7}(\frac{1}{10} + \frac{1}{5})} = 6.6667 \times 10^{-3}$
Iteration 1: $(i = 1)$
 ${}^{2}u^{(1)} = {}^{2}u^{(0)} + \Delta u^{(1)} = 1.3333 \times 10^{-2}$
 ${}^{2}\varepsilon_{a}^{(1)} = 1.3333 \times 10^{-3} < \varepsilon_{Y}$ (elastic section!)
 ${}^{2}\varepsilon_{b}^{(1)} = 2.6667 \times 10^{-3} > \varepsilon_{Y}$ (plastic section!)
 ${}^{1}F_{a}^{(1)} = 1.3333 \times 10^{4};$ ${}^{1}F_{b}^{(1)} = (E^{T}({}^{2}\varepsilon_{b}^{(1)} - \varepsilon_{Y}) + \sigma_{Y})A = 2.0067 \times 10^{4}$
 $({}^{1}K_{a} + {}^{1}K_{b})\Delta u^{(2)} = {}^{2}R - {}^{2}F_{a}^{(1)} - {}^{2}F_{b}^{(1)} \Rightarrow \Delta u^{(2)} = 2.2 \times 10^{-3}$



• We look at the example again – simple bar

i	$\Delta u^{(i)}$	² u ⁽ⁱ⁾
2	1.45E-03	1.55E-02
3	1.45E-03	1.70E-02
4	9.58E-04	1.79E-02
5	6.32E-04	1.86E-02
6	4.17E-04	1.90E-02
7	2.76E-04	1.93E-02





• The basic problem:

We want to establish the solution using an incremental formulation

The equilibrium must be established for the considered body in its current configuration

In proceeding we adopt a Lagrangian formulation where we track the movement of all particles of the body (located in a Cartesian coordinate system)

Another approach would be an Eulerian formulation where the motion of material through a stationary control volume is considered



• The basic problem:





The Lagrangian formulation

We express equilibrium of the body at time $t+\Delta t$ using the principle of virtual displacements

$$\int_{t+\Delta t_V} t^{t+\Delta t} \tau \delta_{t+\Delta t} e_{ij} d^{t+\Delta t} V = t^{t+\Delta t} R$$

$$\delta \mathbf{u} = \begin{bmatrix} \delta u_2 \\ \delta u_3 \end{bmatrix}$$
Configuration corresponding to variation in
displacements $\delta \mathbf{u}$ at ^{*t+Δt*}
Surface area ^{*t+Δt*}
Volume ^{*t+Δt*}
Configuration at time *t* + Δ*t*
Surface area ^{*t+Δt*}
Configuration at time *t*
Surface area ^{*t+Δt*}
Surface area ^{*t+*}

 $\left(\delta u_{1} \right)$

 $t^{t+\Delta t}\tau$: Cartesian components of the Cauchy stress tensor

 $\delta_{t+\Delta t} e_{ij} = \frac{1}{2} \left(\frac{\partial \delta u_i}{\partial^{t+\Delta t} x_j} + \frac{\partial \delta u_j}{\partial^{t+\Delta t} x_i} \right) = \text{strain tensor corresponding to virtual displacements}$

 δu_i : Components of virtual displacement vector imposed at time $t + \Delta t$

 $x_{i}^{t+\Delta t}$: Cartesian coordinate at time $t + \Delta t$

 $^{t+\Delta t}V$: Volume at time $t + \Delta t$

$${}^{t+\Delta t}R = \int_{t+\Delta t_V} {}^{t+\Delta t}f_i^B \delta u_i d^{t+\Delta t}V + \int_{t+\Delta t_{S_f}} {}^{t+\Delta t}f_i^S \delta u_i^S d^{t+\Delta t}S$$



The Lagrangian formulation

We express equilibrium of the body at time $t+\Delta t$ using the principle of virtual displacements

The Lagrangian formulation
We express equilibrium of the body
at time
$$t+\Delta t$$
 using the principle of
virtual displacements
 $t+\Delta t R = \int_{t+\Delta t_V} t+\Delta t \int_{t} B \delta u_i d^{t+\Delta t}V + \int_{t+\Delta t_S_f} t+\Delta t \int_{t} S \delta u_i^S d^{t+\Delta t} S x_1(or^{0}x_1, tx_1, t+\Delta t_N)$

 $f_{i}^{+\Delta t} f_{i}^{B}$: externally applied forces per unit volume $f_{i}^{L+\Delta t}$ f_{i}^{S} : externally applied surface tractions per unit surface $^{t+\Delta t}S_{f}$: surface at time $t + \Delta t$ δu_i^S : δu_i evaluated at the surface ${}^{t+\Delta t}S_f$



• The Lagrangian formulation

We recognize that our derivations from linear finite element theory are unchanged – but applied to the body in the configuration at time $t+\Delta t$



Swiss Federal Institute of Technology

The continuum mechanics incremental equations

• In the further we introduce an appropriate notation:

Coordinates and displacements are related as:

$${}^{t}x_{i} = {}^{0}x_{i} + {}^{t}u_{i}$$
$${}^{t+\Delta t}x_{i} = {}^{0}x_{i} + {}^{t+\Delta t}u_{i}$$

Enti

Increments in displacements are related as:

$$_{t}u_{i}={}^{t+\Delta t}u_{i}-{}^{t}u_{i}$$

Reference configurations are indexed as e.g.:

 ${}^{t+\Delta t}_{0}f_{i}^{S}$ where the lower left index indicates the reference configuration

$${}^{t+\Delta t}\tau_{ij} = {}^{t+\Delta t}_{t+\Delta t}\tau_{ij}$$

Differentiation is indexed as:

$${}^{t+\Delta t}_{0}u_{i,j} = \frac{\partial^{t+\Delta t}u_i}{\partial^{0}x_j}, \qquad {}^{0}_{t+\Delta t}x_{m,n} = \frac{\partial^{0}x_m}{\partial^{t+\Delta t}x_n}$$