## The Finite Element Method for the Analysis of Non-Linear and Dynamic Systems



Prof. Dr. Michael Havbro Faber
Swiss Federal Institute of Technology ETH Zurich, Switzerland


## Contents of Today's Lecture

- Motivation, overview and organization of the course
- Introduction to non-linear analysis
- Formulation of the continuum mechanics incremental equations of motion


## Motivation, overview and organization of the course

- Motivation

In FEM 1 we learned about the steady state analysis of linear systems
however;
the systems we are dealing with in structural engineering are generally not steady state and also not linear

We must be able to assess the need for a particular type of analysis and we must be able to perform it

## Motivation, overview and organization of the course

- Motivation

What kind of problems are not steady state and linear?
E.g. when the:
material behaves non-linearly
deformations become big (p- $\Delta$ effects)
loads vary fast compared to the eigenfrequencies of the structure

General feature: Response becomes load path dependent

## Motivation, overview and organization of the course

- Motivation

What is the "added value" of being able to assess the non-linear non-steady state response of structures?
E.g. assessing the

- structural response of structures to extreme events (rock-fall, earthquake, hurricanes)
- performance (failures and deformations) of soils
- verifying simple models


## Motivation, overview and organization of the course

Steady state problems (Linear/Non-linear):
The response of the system does not change over time

## $\mathbf{K U}=\mathbf{R}$

Propagation problems (Linear/Non-linear):
The response of the system changes over time

$$
\mathbf{M} \ddot{\mathbf{U}}(t)+\mathbf{C} \dot{\mathbf{U}}(t)+\mathbf{K} \mathbf{U}(t)=\mathbf{R}(t)
$$

Eigenvalue problems:
No unique solution to the response of the system

$$
\mathbf{A} \mathbf{v}=\lambda \mathbf{B} \mathbf{v}
$$

## Motivation, overview and organization of the course

- Organisation

The lectures will be given by:

M. H. Faber

Exercises will be organized/attended by:

## J. Qin

By appointment, HIL E13.1

## Motivation, overview and organization of the course

- Organisation

PowerPoint files with the presentations will be uploaded on our homepage one day in advance of the lectures
http://www.ibk.ethz.ch/fa/education/FE_II

The lecture as such will follow the book:
"Finite Element Procedures" by K.J. Bathe, Prentice Hall, 1996

## Motivation, overview and organization of the course

- Overview

| Date | Pages | Subject |
| :---: | :---: | :---: |
| 19.09.2008 | 485-502 | Non-linear Finite Element Calculations in solids and structural mechanics <br> - Introduction to non-linear calculations <br> - The incremental approach to continuum mechanics |
| 26.09.2008 | 502-528 | Non-linear Finite Element Calculations in solids and structural mechanics Deformation gradients, strain and stress tensors The Langrangian formulation - only material non-linearity |
| 03.10.2008 | 538-548 | Non-linear Finite Element Calculations in solids and structural mechanics Displacement based iso-parametric finite elements in continuum mechanics |
| 10.10.2008 | 548-560 | Non-linear Finite Element Calculations in solids and structural mechanics Displacement based iso-parametric finite elements in continuum mechanics |

## Motivation, overview and organization of the course

- Overview

| 17.10.2008 | 561-578 | Non-linear Finite Element Calculations in solids and structural mechanics <br> - Total Langrangian formulation <br> - Extended Langrangian formulation <br> - Structural elements |
| :---: | :---: | :---: |
| 24.10.2008 | 581-617 | Non-linear Finite Element Calculations in solids and structural mechanics Introduction to constitutive relations <br> Non-linear constitutive relations |
| 31.10.2008 | 622-640 | Non-linear Finite Element Calculations in solids and structural mechanics Contact problems <br> Practical considerations |
| 07.11.2008 | 768-784 | $\begin{array}{ll}\text { Dynamical Finite Element Calculations } \\ - & \text { Introduction } \\ - & \text { Direct integration methods }\end{array}$ |



## Motivation, overview and organization of the course

- Overview

| 14.11.2008 | 785-800 | Dynamical Finite Element Calculations - Mode superposition |
| :---: | :---: | :---: |
| 21.11.2008 | 801-815 | Dynamical Finite Element Calculations |
|  |  | Analysis of direct integration methods |
| 28.11.2008 | 824-830 | Dynamical Finite Element Calculations |
|  |  | Solution of dynamical non-linear problems |
| 05.12.2008 | 887-910 | Solution of Eigen value problems |
|  |  | The vector iteration method |
| 12.12.2008 | 911-937 | Solution of Eigen value problems |
|  |  | The transformation method |
| 19.12.2008 |  | Introduction to FEM-software |

## Introduction to non-linear analysis

- Previously we considered the solution of the following linear and static problem:

$$
\mathbf{K U}=\mathbf{R}
$$

for these problems we have the convenient property of linearity, i.e.:

$$
\begin{aligned}
& \mathbf{K} \mathbf{U}^{*}=\lambda \mathbf{R} \\
& \Downarrow \\
& \mathbf{U}^{*}=\lambda \mathbf{U}
\end{aligned}
$$

If this is not the case we are dealing with a non-linear problem!

## Introduction to non-linear analysis

- Previously we considered the solution of the following linear and static problem:
$\mathbf{K U}=\mathbf{R}$
we assumed:
small displacements when developing the stiffness matrix K and the load vector $R$, because we performed all integrations over the original element volume
that the B matrix is constant independent of element displacements
the stress-strain matrix $C$ is constant
boundary constraints are constant



## Introduction to non-linear analysis

## - Classification of non-linear analysis

| Type of analysis | Description | Typical <br> formulation used | Stress and strain <br> measures used |
| :--- | :--- | :--- | :--- |
| Materially-nonlinear <br> only | Infinitesimal <br> displacements and <br> strains; stress train <br> relation is non- <br> linear | Materially- <br> nonlinear-only <br> (MNO) | Engineering strain <br> and stress |
| Large <br> displacements, large <br> rotations but small <br> strains | Displacements and <br> rotations of fibers <br> are large; but fiber <br> extensions and <br> angle changes <br> between fibers are <br> small; stress strain <br> relationship may be <br> linear or non-linear | Total Lagrange (TL) <br> (UL) | Second Piola- <br> Kirchoff stress, <br> Green-Lagrange <br> strain |
| Large <br> displacements, large <br> rotations and large <br> strains | Displacements and <br> rotations of fibers <br> are large; fiber <br> extensions and <br> angle changes <br> between fibers may <br> also be large; stress <br> strain relationship <br> may be linear or <br> non-linear | Total Lagrange (TL) | Cauchy stress, <br> Almansi strain |
| (UL) |  |  |  |$\quad$| Second Piola- |
| :--- |
| Kirchoff stress, |
| Green-Lagrange |
| strain |\(\left|-\begin{array}{l}Cauchy stress, <br>

Logarithmic strain\end{array}\right|\)

## Introduction to non-linear analysis

- Classification of non-linear analysis



Linear elastic (infinitesimal displacements)


## Introduction to non-linear analysis

- Classification of non-linear analysis



## Introduction to non-linear analysis



Large displacements and large rotations but small strains (linear or nonlinear material behavior)


## Introduction to non-linear analysis

- Classification of non-linear analysis


Large displacements and large rotations and large strains (linear or nonlinear material behavior)


## Introduction to non-linear analysis

- Classification of non-linear analysis


Changing boundary conditions

## Introduction to non-linear analysis

- Example: Simple bar structure




$$
\varepsilon_{Y}=0.002
$$



## Introduction to non-linear analysis

- Example: Simple bar structure




$$
\begin{aligned}
& { }^{t} \varepsilon_{a}=\frac{{ }^{t} u}{L_{a}},{ }^{t} \varepsilon_{b}=-\frac{{ }^{t} u}{L_{b}} \\
& { }^{t} R+{ }^{t} \sigma_{b} A={ }^{t} \sigma_{a} A \\
& { }^{t} \varepsilon=\frac{{ }^{t} \sigma}{E} \text { (elastic region) }
\end{aligned}
$$

$$
{ }^{t} \varepsilon=\varepsilon_{Y}+\frac{{ }^{t} \sigma-\sigma_{Y}}{E_{T}} \text { (plastic region) } \quad \Delta \varepsilon=\frac{\Delta \sigma}{E} \text { (unloading) }
$$

## Introduction to non-linear analysis

- Example: Simple bar structure


$$
\begin{aligned}
& { }^{t} \varepsilon_{a}=\frac{{ }^{t} u}{L_{a}},{ }^{t} \varepsilon_{b}=-\frac{{ }^{t} u}{L_{b}} \\
& { }^{t} R+{ }^{t} \sigma_{b} A={ }^{t} \sigma_{a} A \\
& { }^{t} \varepsilon=\frac{{ }^{t} \sigma}{E} \text { (elastic region) } \\
& { }^{t} \varepsilon=\varepsilon_{Y}+\frac{{ }^{t} \sigma-\sigma_{Y}}{E_{T}} \text { (plastic region) } \\
& \Delta \varepsilon=\frac{\Delta \sigma}{E} \text { (unloading) }
\end{aligned}
$$

## Both sections elastic

$$
\begin{aligned}
& { }^{t} R=E A^{t} u\left(\frac{1}{L_{a}}+\frac{1}{L_{b}}\right) \Rightarrow{ }^{t} u=\frac{{ }^{t} R}{3 \cdot 10^{6}} \\
& \sigma_{a}=\frac{{ }^{t} R}{3 A}, \sigma_{b}=-\frac{2}{3} \frac{{ }^{t} R}{A}
\end{aligned}
$$



## Introduction to non-linear analysis

- Example: Simple bar structure


$$
\begin{aligned}
& { }^{t} \varepsilon_{a}=\frac{{ }^{t} u}{L_{a}}, \varepsilon_{b}=-\frac{{ }^{t} u}{L_{b}} \\
& { }^{t} R+{ }^{t} \sigma_{b} A={ }^{t} \sigma_{a} A \\
& { }^{t} \mathcal{E}=\frac{{ }^{t} \sigma}{E} \text { (elastic region) } \\
& { }^{t} \mathcal{E}=\varepsilon_{Y}+\frac{{ }^{t} \sigma-\sigma_{Y}}{E_{T}} \text { (plastic region) } \\
& \Delta \varepsilon=\frac{\Delta \sigma}{E} \text { (unloading) }
\end{aligned}
$$

## Section $\mathbf{a}$ is elastic while section $\mathbf{b}$ is plastic

section ${ }_{t_{u}}$ will be plastic when ${ }^{t_{u}} R=\frac{3}{2} \sigma_{Y} A$
$\sigma_{a}=E \frac{{ }^{t} u}{L_{a}}, \sigma_{b}=E_{T}\left(\frac{{ }^{t} u}{L_{b}}-\varepsilon_{Y}\right)-\sigma_{Y}$
${ }^{t} R=\frac{E A^{t} u}{L_{a}}+\frac{E_{T} A^{t} u}{L_{b}}-E_{T} \varepsilon_{Y} A+\sigma_{Y} A \Rightarrow$
${ }^{t} u=\frac{{ }^{t} R / A+E_{T} \varepsilon_{Y}-\sigma_{Y}}{E / L_{a}+E / L_{b}}=\frac{{ }^{t} R}{1.02 \cdot 10^{6}}-1.9412 \cdot 10^{-2}$

## Introduction to non-linear analysis

- What did we learn from the example?

The basic problem in general nonlinear analysis is to find a state of equilibrium between externally applied loads and element nodal forces

$$
\begin{aligned}
& { }^{t} \mathbf{R}-{ }^{t} \mathbf{F}=0 \\
& { }^{t} \mathbf{R}={ }^{t} \mathbf{R}_{B}+{ }^{t} \mathbf{R}_{S}+{ }^{t} \mathbf{R}_{C} \\
& { }^{t} \mathbf{F}={ }^{t} \mathbf{R}_{I} \\
& { }^{t} \mathbf{F}=\sum_{m} \int_{V^{(m)}}{ }^{t} \mathbf{B}^{(m) T} \tau^{(m) t} d V^{(m)}
\end{aligned}
$$

We must achieve equilibrium for all time steps when incrementing the loading

Very general approach
includes implicitly also dynamic analysis!

## Introduction to non-linear analysis

- The basic approach in incremental anaylsis is
${ }^{t+\Delta t} \mathbf{R}-{ }^{t+\Delta t} \mathbf{F}=0$
assuming that ${ }^{t+\Delta t} \mathbf{R}$ is independent of the deformations we have
${ }^{t+\Delta t} \mathbf{F}={ }^{t} \mathbf{F}+\mathbf{F}$
We know the solution ${ }^{t} F$ at time $t$ and $F$ is the increment in the nodal point forces corresponding to an increment in the displacements and stresses from time $t$ to time $t+\Delta t$ this we can approximate by
$\mathbf{F}={ }^{t} \mathbf{K} \mathbf{U}$
$\dagger$
Tangent stiffness matrix $\quad{ }^{t} \mathbf{K}=\frac{\partial^{t} \mathbf{F}}{\partial^{t} \mathbf{U}}$


## Introduction to non-linear analysis

- The basic approach in incremental anaylsis is

We may now substitute the tangent stiffness matrix into the equlibrium relation

```
't}\mathbf{KU}=\mp@subsup{}{}{t+\Deltat}\mathbf{R}-\mp@subsup{}{}{t}\mathbf{F
\Downarrow
*+\Deltat}\mathbf{U}=\mp@subsup{}{}{t}\mathbf{U}+\mathbf{U
```

which gives us a scheme for the calculation of the displacements
the exact displacements at time $t+\Delta t$ correspond to the applied loads at $\boldsymbol{t}+\Delta \boldsymbol{t}$ however we only determined these approximately as we used a tangent stiffness matrix - thus we may have to iterate to find the solution


## Introduction to non-linear analysis

- The basic approach in incremental anaylsis is

We may use the Newton-Raphson iteration scheme to find the equlibrium within each load increment

$$
\begin{aligned}
& { }^{t+\Delta t} \mathbf{K}^{(i-1)} \Delta \mathbf{U}^{(i)}={ }^{t+\Delta t} \mathbf{R}-{ }^{t+\Delta t} \mathbf{F}^{(i-1)} \quad \text { (out of balance load vector) } \\
& { }^{t+\Delta t} \mathbf{U}^{(i)}={ }^{t+\Delta t} \mathbf{U}^{(i-1)}+\Delta \mathbf{U}^{(i)} \\
& \text { with initial conditions } \\
& { }^{t+\Delta t} \mathbf{U}^{(0)}={ }^{t} \mathbf{U} ; \quad{ }^{t+\Delta t} \mathbf{K}^{(0)}={ }^{t} \mathbf{K} ; \quad{ }^{t+\Delta t} \mathbf{F}^{(0)}={ }^{t} \mathbf{F}
\end{aligned}
$$

## Introduction to non-linear analysis

- The basic approach to incremental analysis is

It may be expensive to calculate the tangent stiffness matrix and;
in the Modified Newton-Raphson iteration scheme it is thus only calculated in the beginning of each new load step
in the quasi-Newton iteration schemes the secant stiffness matrix is used instead of the tangent matrix

## The continuum mechanics incremental equations

- The basic problem:

We want to establish the solution using an incremental formulation

The equilibrium must be established for the considered body in ist current configuration

In proceeding we adopt a Lagrangian formulation where track the movement of all particles of the body (located in a Cartesian coordinate system)

Another approach would be an Eulerian formulation where the motion of material through a stationary control volume is considered

## The continuum mechanics incremental equations

- The basic problem:



## The continuum mechanics incremental equations

- The Lagrangian formulation

We express equilibrium of the body at time $\boldsymbol{t + \Delta t}$ using the principle of virtual displacements
$\int_{{ }^{+}+\Delta_{V}}{ }^{t+\Delta t} \tau \delta_{t+\Delta t} e_{i j} d^{t+\Delta t} V={ }^{t+\Delta t} R$

${ }^{t+\Delta t} \tau$ : Cartesian components of the Cauchy stress tensor
$\delta_{t+\Delta t} e_{i j}=\frac{1}{2}\left(\frac{\partial \delta u_{i}}{\partial^{t+\Delta t} x_{j}}+\frac{\partial \delta u_{j}}{\partial^{t+\Delta t} x_{i}}\right)=$ strain tensor corresponding to virtual displacements
$\delta u_{i}$ : Components of virtual displacement vector imposed at time $t+\Delta t$
${ }^{t+\Delta t} x_{i}$ : Cartesian coordinate at time $t+\Delta t$
${ }^{t+\Delta t} V$ : Volume at time $t+\Delta t$
${ }^{t+\Delta t} R=\int_{{ }^{t+\Delta^{t}} V}{ }^{t+\Delta t} f_{i}^{B} \delta u_{i} d^{t+\Delta t} V=\int_{t+\Delta S_{S}}{ }^{t+\Delta t} f_{i}^{S} \delta u_{i}^{S} d^{t+\Delta t} S$


## The continuum mechanics incremental equations

- The Lagrangian formulation

We express equilibrium of the body at time $\boldsymbol{t + \Delta t}$ using the principle of virtual displacements


where
${ }^{t+\Delta t} f_{i}^{B}$ : externally applied forces per unit volume
${ }^{t+\Delta t} f_{i}^{S}$ : externally applied surface tractions per unit surface
${ }^{t+\Delta t} S_{f}$ : surface at time $t+\Delta t$
$\delta u_{i}^{S}: \delta u_{i}$ evaluated at the surface ${ }^{t+\Delta t} S_{f}$

## The continuum mechanics incremental equations

- The Lagrangian formulation

We recognize that our derivations from linear finite element theory are unchanged - but applied to the body in the configuration at time $\boldsymbol{t}+\Delta \boldsymbol{t}$

## The continuum mechanics incremental equations

- In the further we introduce an appropriate notation:

Coordinates and displacements are related as:
${ }^{t} x_{i}={ }^{0} x_{i}+{ }^{t} u_{i}$
${ }^{t+\Delta t} x_{i}={ }^{0} x_{i}+{ }^{t+\Delta t} u_{i}$
Increments in displacements are related as:
${ }_{t} u_{i}={ }^{t+\Delta t} u_{i}-{ }^{t} u_{i}$
Reference configurations are indexed as e.g.:
${ }^{t+\Delta t}{ }_{0} f_{i}^{S}$ where the lower left index indicates the reference configuration

$$
{ }^{t+\Delta t} \tau_{i j}={ }_{t+\Delta t}^{t+\Delta t} \tau_{i j}
$$

Differentiation is indexed as:

$$
{ }_{0}^{t+\Delta t} u_{i, j}=\frac{\partial^{t+\Delta t} u_{i}}{\partial^{0} x_{j}}, \quad{ }_{t+\Delta t}^{0} x_{m, n}=\frac{\partial^{0} x_{m}}{\partial^{t+\Delta t} x_{n}}
$$

